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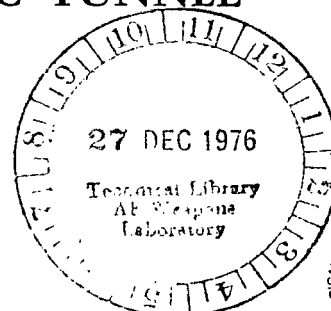


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## INVESTIGATION OF VERY LOW BLOCKAGE RATIO BOATTAIL MODELS IN THE LANGLEY 16-FOOT TRANSONIC TUNNEL

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# INVESTIGATION OF VERY LOW BLOCKAGE RATIO BOATTAIL MODELS IN THE LANGLEY 16-FOOT TRANSONIC TUNNEL

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## SUMMARY

An investigation at an angle of attack of  $0^\circ$  has been conducted in the Langley 16-foot transonic tunnel at Mach numbers from 0.4 to 1.05 to determine the limits in Mach number at which valid boattail pressure drag data may be obtained with very low blockage (0.0027) ratio bodies. The results of this investigation indicate that extreme care must be exercised when examining any data taken at subsonic Mach numbers very near 1.0 (e.g., above  $M = 0.96$ ) and lower than the supersonic Mach number at which shock reflections miss the model. Boattail pressure coefficient distributions will generally not indicate any error, but when integrated boattail pressure drag data are plotted as a function of Mach number, data which may be in error can generally be identified.

## INTRODUCTION

Even with slotted or porous wall wind tunnels, the difficulties in obtaining valid wind-tunnel data near a Mach number of 1.0 have long been known. At Mach numbers very slightly less than 1.0, although the models used may have relatively small blockage, regions of sonic flow will extend to the wind-tunnel walls. The shape of the sonic regions is changed by the interaction with the walls and the resulting data are then in error. At Mach numbers slightly greater than 1.0, the model bow shock is reflected back from the walls and impinges on the model, again resulting in invalid data. An investigation was recently conducted in the Langley 0.3-meter transonic cryogenic tunnel to determine the effects of Reynolds number on boattail pressure drag (ref. 1). In order to accomplish the objectives of this investigation, six small models (2.54 cm in diameter) with extensive boattail pressure instrumentation were constructed. With these small models available, an investigation was then initiated in the Langley 16-foot transonic tunnel to determine the practical testing limit for Mach number (near 1.0) at which valid boattail pressure data could be obtained since the models would have an extremely low blockage ratio (ratio of model cross-sectional area to tunnel test section cross-sectional area) in this facility. In addition, a facility-to-facility comparison between the 0.3-meter transonic cryogenic

tunnel and the 16-foot transonic tunnel could be made. Also, two of the small models were scale models of two large models which had been previously tested in the 16-foot transonic tunnel (refs. 2 and 3), and an additional data comparison of the effect of model size in the same facility would be possible.

The current investigation used the series of six isolated, sting-mounted, cone-cylinder nacelle models with four different boattail geometries (ref. 1). These models were tested in the Langley 16-foot transonic tunnel at  $0^\circ$  angle of attack at Mach numbers from about 0.4 to 1.05. Reynolds number based on the distance from the nose of the model to the start of the boattail varied from about  $1.7 \times 10^6$  at  $M = 0.4$  to  $2.9 \times 10^6$  at  $M = 1.05$  for four of the models and from about  $3.4 \times 10^6$  at  $M = 0.4$  to  $5.7 \times 10^6$  at  $M = 1.05$  for the other two models.

## SYMBOLS

$A$	cross-sectional area
$A_m$	maximum cross-sectional area of model
$A_\beta$	incremental cross-sectional area assigned to boattail static-pressure orifice for drag integration
$C_{D,\beta}$	boattail pressure drag coefficient (see Data Reduction section)
$C_p$	static-pressure coefficient, $\frac{p - p_\infty}{q}$
$d_m$	maximum diameter of model
$d_s$	sting diameter
$L$	length of model from nose to start of boattail (characteristic length)
$l$	length of boattail
$M$	free-stream Mach number
$p$	local static pressure on model
$p_\infty$	free-stream static pressure

$q$	free-stream dynamic pressure
$R$	Reynolds number (based on length from nose to start of boattail)
$x$	axial distance from start of boattail, positive aft
$y$	radial distance from center line of model
$\phi$	meridian angle about model axis, clockwise positive facing upstream, $0^\circ$ at top of model

## APPARATUS AND PROCEDURE

### Wind Tunnel

This investigation was conducted in the Langley 16-foot transonic tunnel, which is a single-return, continuous-flow, atmospheric tunnel. The test section is a regular octagon in cross section with slots at the corners of the octagon. The tunnel speed can be varied continuously from a Mach number of 0.20 to 1.30. Further description of the Langley 16-foot transonic tunnel can be found in references 4 to 6.

### Models and Support System

A generalized sketch of the boattailed cone-cylinder nacelle models used in this investigation is shown in figure 1. Figure 2 is a photograph of all six models. There were four short models of differing boattail geometry with a length of 20.32 cm from the nose to the start of the boattail (characteristic length); there were two long models with a length from the nose to the start of the boattail of 40.64 cm. The boattail geometry of the two long models duplicated the boattail geometry of two of the short models. Details of the geometry of the four boattails are shown in figure 3. The four boattail geometries were a circular arc with a ratio of length to maximum diameter or fineness ratio  $l/d_m$  of 0.8 (both short and long models), circular arc with a fineness ratio of 1.77, circular-arc-conic with a fineness ratio of 0.96 (both short and long models), and contoured with a fineness ratio of 0.95. The two circular-arc boattails are scale models of two boattails which have been tested in the Langley 16-foot transonic tunnel (refs. 2 and 3).

The models were all sting mounted with the sting simulating the geometry of a jet exhaust plume for a nozzle operating at its design point. (See ref. 3.) The two circular-arc boattails and the circular-arc-conic boattail had ratios of sting diameter to maximum diameter of 0.50, whereas the contoured nozzle had a ratio of sting diameter to maximum diameter of 0.544. The length of the constant-diameter portion of the stings

was such that, based on the data contained in reference 7, there should be no effect of the tunnel support sting flare on the boattail pressure coefficients. In addition, the sum of the boattail and sting lengths (before the flare) was constant so that the noses of all four of the short models were at the same tunnel station. The noses of the two long models were at the same tunnel station, but this station was not the same as that for the short models. (The start of the boattail was at the same tunnel station for all six models.)

The models were constructed of cast aluminum with stainless-steel pressure tubes cast as an integral part of each model. The tubes were placed in the sand mold in the proper position, the aluminum was poured, and the model was machined to the proper contours.

Figure 4 shows photographs of a typical short model and a typical long model mounted in the Langley 16-foot transonic tunnel. The models had a maximum diameter of 2.54 cm and a resulting tunnel blockage of 0.0027 percent.

#### Instrumentation and Tests

The six boattails were instrumented with 30 static-pressure orifices in 3 rows of 10 orifices each ( $\phi = 0^\circ, 120^\circ, \text{ and } 240^\circ$ ) at the locations given in table I. These orifices were connected to remotely located 34.47 kPa (5 psi) pressure gages which had an accuracy of  $\pm 0.5$  percent of full scale.

All tests were conducted in the Langley 16-foot transonic tunnel at Mach numbers from 0.40 to 1.05 at an angle of attack of  $0^\circ$ . Model attitude was set to account for tunnel upflow (about  $0.1^\circ$  through the Mach number range), but no account was taken of possible sting deflection because it was believed to be insignificant. Boundary-layer transition was natural for all tests to insure a direct comparison with the data from reference 1.

#### DATA REDUCTION

Model and wind-tunnel data were recorded on magnetic tape and a digital computer was used to compute standard force and pressure coefficients. Pressure drag coefficients were computed from the measured pressures on each boattail. These coefficients, based on the maximum cross-sectional area of the model, were obtained from the pressure data on the boattail portion of the model by assigning an axially projected area to each orifice and computing the coefficients from the following equation:

$$C_{D,\beta} = \frac{1}{qA_m} \sum_{i=1}^{30} (p_i - p_\infty) A_{\beta,i}$$

Accuracy of this step integration scheme was spot-checked by plotting the pressure coefficients as a function of  $A/A_m$  and integrating with a planimeter.

## DISCUSSION

### Boattail Pressure Coefficient Distributions

Boattail pressure coefficient distributions for the six models at the various Mach numbers tested are shown in figures 5 to 10. The significant point to be made about these data is that, even though the data obtained above  $M = 1.0$  will later be shown to be in error and some data obtained just below  $M = 1.0$  (e.g., above  $M = 0.96$ ) becomes questionable when the integrated boattail pressure drag coefficients are examined, it is impossible to ascertain such errors by looking only at the boattail pressure coefficient distributions. Figure 11 shows typical examples of boattail pressure coefficient distributions for two of the boattails at the Mach numbers tested between about 0.9 and 1.0. (Both the actual data and a fairing through the data are shown since both the  $\phi = 0^\circ$  and  $\phi = 120^\circ$  rows of pressures are included in the distributions.) Even when boattail pressure coefficient distributions at successive Mach numbers are plotted, the changes in distributions due to the rearward shock movement with Mach number seem reasonable and no error can readily be ascertained.

### Boattail Pressure Drag Coefficients

Boattail pressure drag coefficients as a function of Mach number for the six configurations are shown in figure 12. When these curves are examined, it is immediately apparent that the data above  $M = 1.0$  are in error. The integrated drag coefficients are not in a smooth curve with Mach number and appear to be shifting up and down slightly. At Mach numbers below 1.0, there is also some question as to the validity of the data for tests above  $M = 0.96$ . At about  $M = 0.98$ , the data depart from a smoothly faired curve for all six configurations, and although the data above  $M = 0.98$  may generally be faired into a continuation of the lower Mach number curve, this phenomenon leads one to believe that these data at  $M = 0.98$  and above are possibly in error. This indication that the onset of transonic wall interference occurs at a Mach number of 0.96 agrees with the results found in reference 8, which are also for low blockage ratio bodies.

It is believed that the phenomenon above  $M = 1.0$  is due to the reflected bow shock impinging on the model and would be expected to clear up once a supersonic Mach number is reached at which the reflected shock misses the model. It has been suggested that the phenomenon below  $M = 1.0$  is due to a sonic region on the model extending to the wall and its shape being changed through interaction with the wall, but available analytic procedures (e.g., ref. 9) indicate that the sonic region reaches the wall only at Mach numbers greater than 0.99 (for this geometry and blockage ratio).

## Comparison of Data

Boattail pressure coefficient distributions.- Boattail pressure coefficient distributions for the  $L/d_m = 16.0$ ,  $l/d_m = 0.80$  circular-arc boattail, the  $L/d_m = 8.0$  and  $16.0$ ,  $l/d_m = 0.96$  circular-arc-conic boattails, and the  $L/d_m = 8.0$ ,  $l/d_m = 0.95$  contoured boattail obtained in both the Langley 16-foot transonic tunnel (blockage about 0.0027 percent) and the Langley 0.3-meter transonic cryogenic tunnel (ref. 1, blockage about 0.52 percent) are shown in figures 13 to 16 for  $M = 0.6$  and  $0.9$ . The other two configurations ( $L/d_m = 8.0$ ,  $l/d_m = 0.80$  and  $1.77$  circular-arc boattails) were not tested at a Reynolds number in the cryogenic tunnel low enough to be sufficiently close to the 16-foot transonic tunnel Reynolds number for a valid comparison. It was shown in reference 1 that although integrated boattail pressure drag coefficients are not sensitive to Reynolds number, the boattail pressure coefficient distributions are sensitive to Reynolds number, and thus while comparisons of pressure drag coefficients are valid at different Reynolds numbers, comparisons of pressure coefficient distributions are not. For those four configurations for which comparisons can be made, only a general trend can be ascertained. The boattail pressure coefficient distributions are generally in reasonable agreement between the two tunnels with the  $M = 0.9$  data being in slightly better agreement than the  $M = 0.6$  data.

Boattail pressure drag coefficients in the Langley 16-foot transonic tunnel and the Langley 0.3-meter transonic cryogenic tunnel.- Boattail pressure drag coefficients for the data from both the 16-foot tunnel (blockage about 0.0027 percent) and the cryogenic tunnel (ref. 1, blockage about 0.52 percent) and for the large model data from the 16-foot tunnel (ref. 3, blockage of model and support about 0.15 percent) are shown in figure 17. For all six small models, the boattail pressure drag coefficients obtained in both the 16-foot tunnel and the cryogenic tunnel agree quite well at Mach numbers of 0.8 and above. However, the drag levels for the models in the cryogenic tunnel are always lower than those for the models in the 16-foot tunnel at  $M = 0.6$  since the pressure coefficients are slightly higher for the cryogenic tunnel than for the 16-foot tunnel at  $M = 0.6$ . The reason for such a difference between the drag levels is not clear, but possibly may be due to the fact that the cryogenic tunnel has wider slots than the 16-foot tunnel and hence a greater open area. It is not believed that the difference in blockage causes the difference in drag levels because the Mach number is low and the data agree at the higher Mach numbers.

Boattail pressure drag coefficients for two model sizes in the Langley 16-foot transonic tunnel.- For both the large (refs. 2 and 3) and small versions of the two circular-arc boattails ( $L/d_m = 8.0$ ,  $l/d_m = 0.8$  and  $1.77$ ) which were tested in the 16-foot tunnel, the results are not the same. The drag coefficients for the  $l/d_m = 1.77$  circular-arc boattail which has attached flow at most Mach numbers agree extremely well, whereas



the drag level of the large  $l/d_m = 0.80$  circular-arc boattail which has separated flow over the rear portion of the boattail (as indicated by the flattening out of the pressure distribution to a constant level) at all Mach numbers tested is somewhat higher than that of the small model. It is believed that this discrepancy for the  $l/d_m = 0.80$  boattail occurs because boundary-layer transition was fixed at the nose for the large models while it was natural for the small models. The resulting differences in the boundary layer approaching the boattails for the two different size models resulted in changes in the boattail pressure coefficient distribution for the  $l/d_m = 0.80$  boattail with separated flow, while not for the  $l/d_m = 1.77$  boattail with attached flow. The sensitivity of the separated-flow boattail to changes in approach boundary layer is not unexpected. Note that for the two large models the drag data above  $M = 0.96$  are in error and exhibit a tendency to drop off as  $M = 1.0$  is approached. For this blockage ratio in the 16-foot tunnel (0.15 percent), the data in error are easier to detect than for the small blockage ratio models since the effects are larger.

## CONCLUSIONS

An investigation at an angle of attack of  $0^\circ$  has been conducted in the Langley 16-foot transonic tunnel at Mach numbers from about 0.4 to 1.05 to determine the limits in Mach number at which valid boattail pressure drag data may be obtained for very low blockage (0.0027 percent) ratio bodies.

First, it was found that extreme care must be exercised when examining any boat-tail pressure data taken at subsonic Mach numbers greater than about 0.96 and lower than the supersonic Mach number at which shock reflections miss the model. Boattail pressure coefficient distributions will generally not indicate any error, but when integrated boattail pressure drag data are plotted as a function of Mach number, data which may be in error can generally be identified.

Second, comparisons between boattail pressure drag data obtained using both small and large models in the Langley 16-foot transonic tunnel and the data obtained for the small models tested in the Langley 0.3-meter transonic cryogenic tunnel were generally very good.

Langley Research Center  
National Aeronautics and Space Administration  
Hampton, VA 23665  
September 30, 1976

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TABLE I.- BOATTAIL STATIC-PRESSURE ORIFICE LOCATIONS

Boattail configuration	$\frac{x}{d_m}$ for $\frac{L}{d_m} = 8$ at -			$\frac{x}{d_m}$ for $\frac{L}{d_m} = 16$ at -		
	$\phi = 0^\circ$	$\phi = 120^\circ$	$\phi = 240^\circ$	$\phi = 0^\circ$	$\phi = 120^\circ$	$\phi = 240^\circ$
Circular-arc; $\frac{l}{d_m} = 0.8$	-0.2771	-0.2761	-0.2850	-0.4491	-0.4660	-0.4561
	-.0256	-.0731	-.0700	-.1637	-.2201	-.1552
	.0770	.0256	.0345	-.0600	-.1281	-.0590
	.1765	.1287	.1270	.0337	-.0260	.0390
	.2750	.2257	.2260	.1268	.0744	.1342
	.3679	.3240	.3279	.2279	.1729	.2713
	.4675	.4180	.4200	.3210	.2696	.3718
	.5749	.5166	.5220	.4199	.3679	.4680
	.6698	.6165	.6376	.5231	.4640	.5749
	.7746	.7280	.7400	.6279	.6758	.7304
Circular-arc; $\frac{l}{d_m} = 1.77$	-0.3006	-0.3028	-0.3090			
	-.0086	.2554	-.1088			
	.3515	.5057	-.0113			
	.6314	.7364	.3543			
	.8271	.9091	.5072			
	.9975	1.0652	.8204			
	1.1500	1.2075	.9910			
	1.2776	1.3581	1.3528			
	1.4144	1.4817	1.4846			
	1.5447	1.6202	1.6936			
Circular-arc-conic; $\frac{l}{d_m} = 0.96$	-0.2842	-0.2843	-0.2847	-0.2612	-0.2773	-0.2613
	.0001	-.0352	.0059	.0228	-.0292	.0307
	.1090	.0640	.1044	.1227	.0707	.1246
	.1917	.1577	.2069	.2158	.1705	.2308
	.3000	.2542	.3031	.3107	.2638	.3237
	.4027	.3528	.4453	.4148	.3638	.4648
	.5022	.4499	.5483	.5097	.4607	.5737
	.5991	.5518	.6516	.5967	.5557	.6628
	.6948	.6522	.7494	.7068	.6558	.7706
	.8036	.8488	.9069	.8188	.8678	.9156
Contoured; $\frac{l}{d_m} = 0.95$	-0.2951	-0.2950	-0.3011			
	-.0084	-.0500	-.0021			
	.0925	.0500	.0934			
	.1930	.1485	.1976			
	.2861	.2433	.2913			
	.3759	.3449	.4350			
	.4786	.4417	.5355			
	.5797	.5404	.6335			
	.6840	.6415	.7366			
	.7909	.8425	.8951			

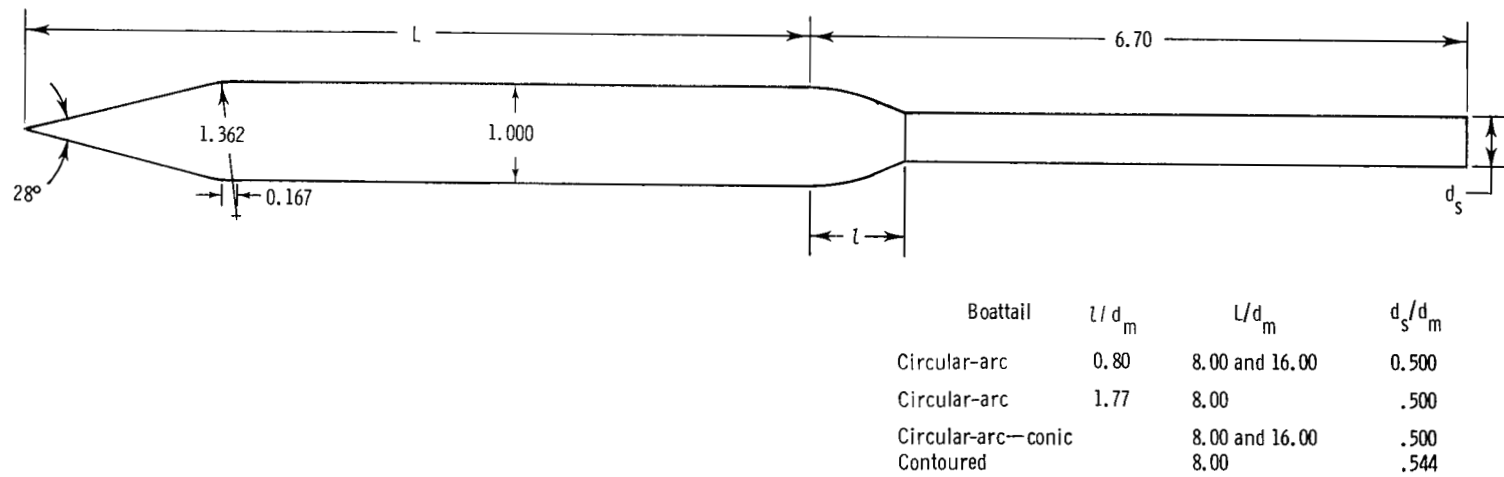
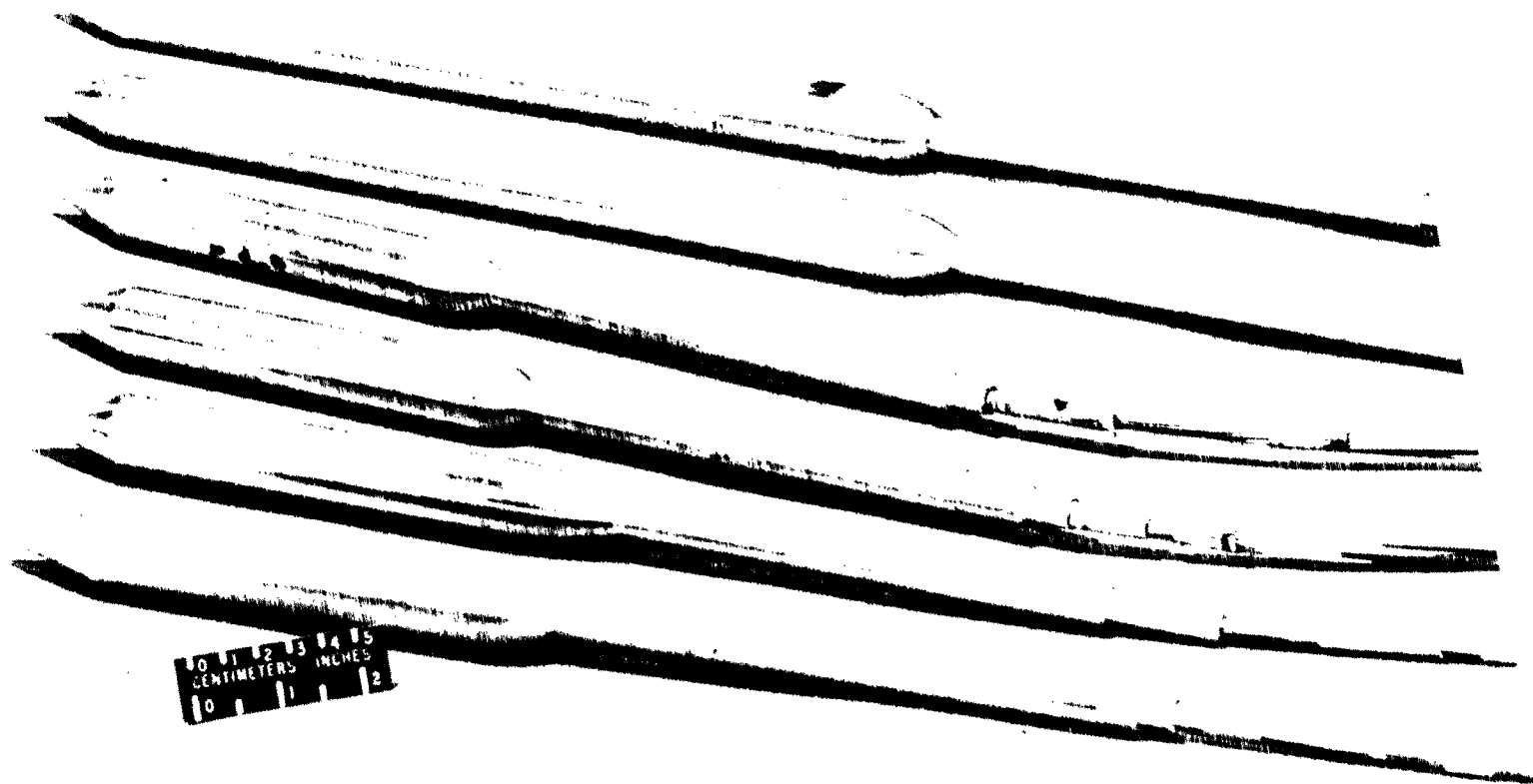


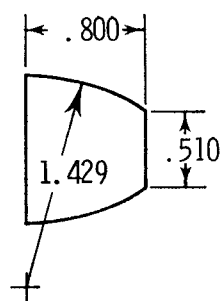
Figure 1.- Boattailed cone-cylinder nacelle model. All dimensions are nondimensionalized by model maximum diameter (2.54 cm).



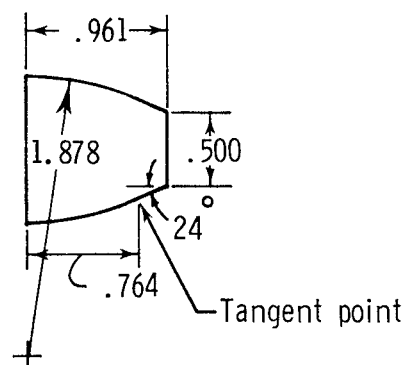
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Figure 2.- Six nacelle models used in this investigation.

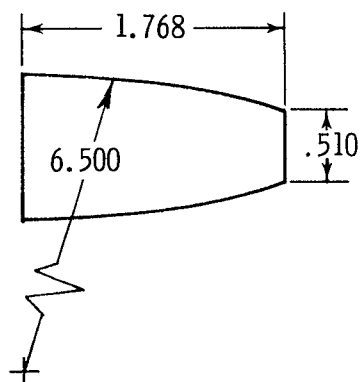
Circular-arc



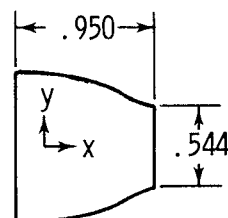
Circular-arc—conic



Circular-arc



Contoured



Coordinates for contoured boattail

$x/d_m$	$y/d_m$	$x/d_m$	$y/d_m$
0.0	0.500	0.640	0.409
.080	.499	.720	.372
.160	.497	.800	.328
.240	.494	.840	.304
.320	.487	.880	.288
.400	.475	.920	.276
.480	.461	.950	.272
.560	.439		

Figure 3.- Details of boattail geometries. All dimensions are nondimensionalized by model maximum diameter (2.54 cm).



L-74-8664

(a) Front view of typical  $L/d_m = 8.0$  model.

Figure 4.- Typical model installed in Langley 16-foot transonic tunnel.



(b) Side view of typical  $L/d_m = 8.0$  model.

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Figure 4.- Continued.





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(c) Front view of typical  $L/d_m = 16.0$  model.

Figure 4.- Concluded.

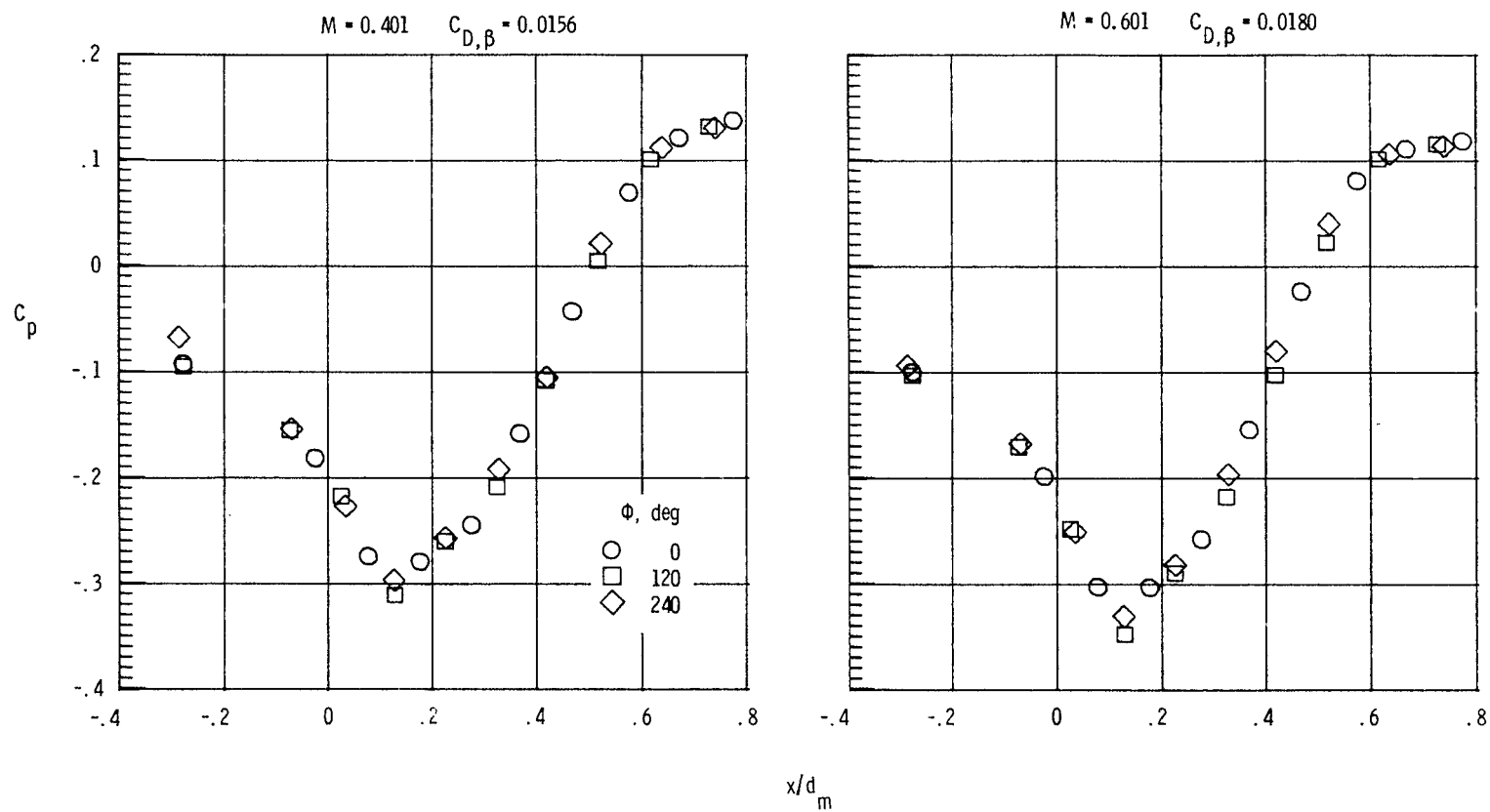
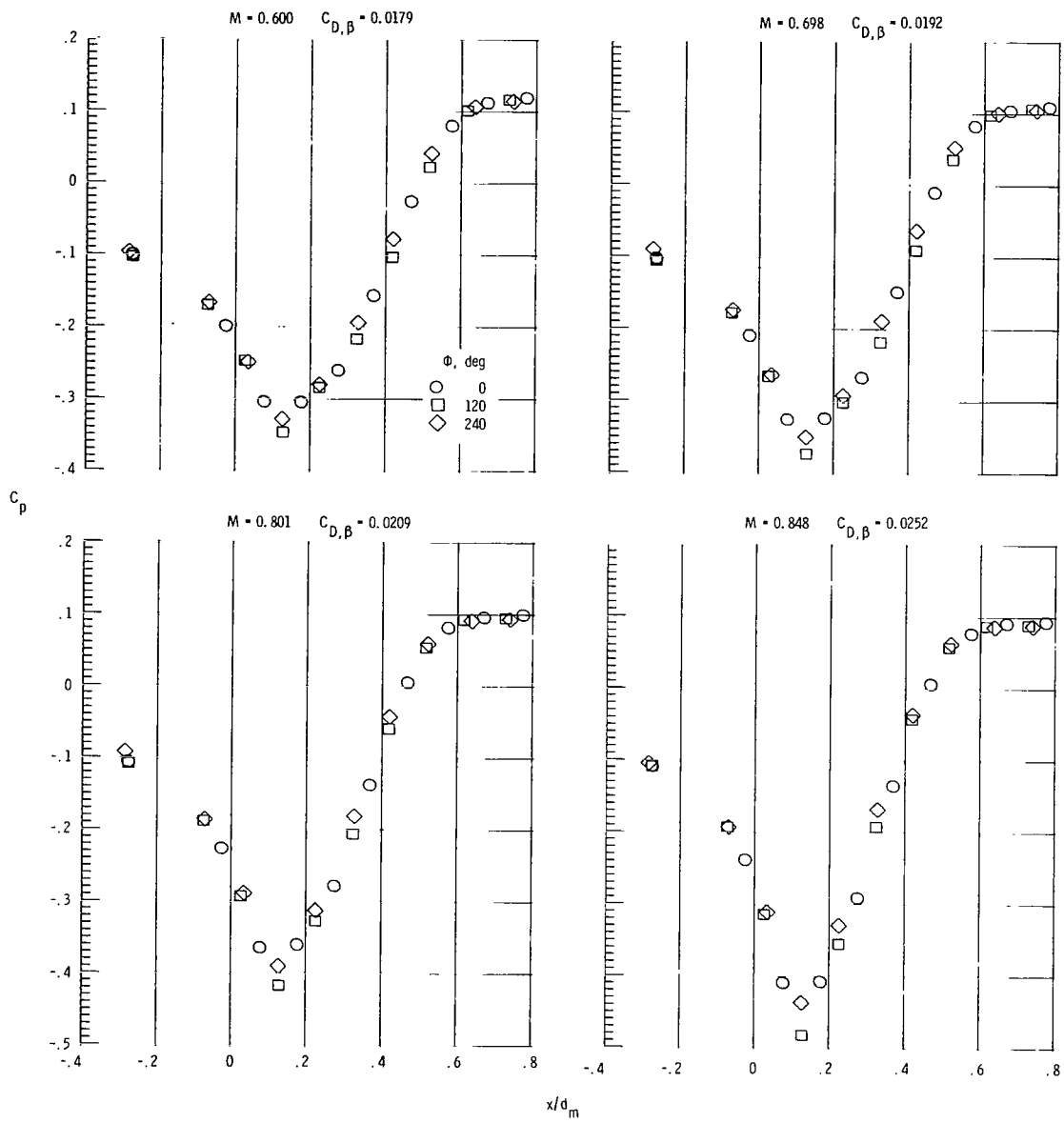
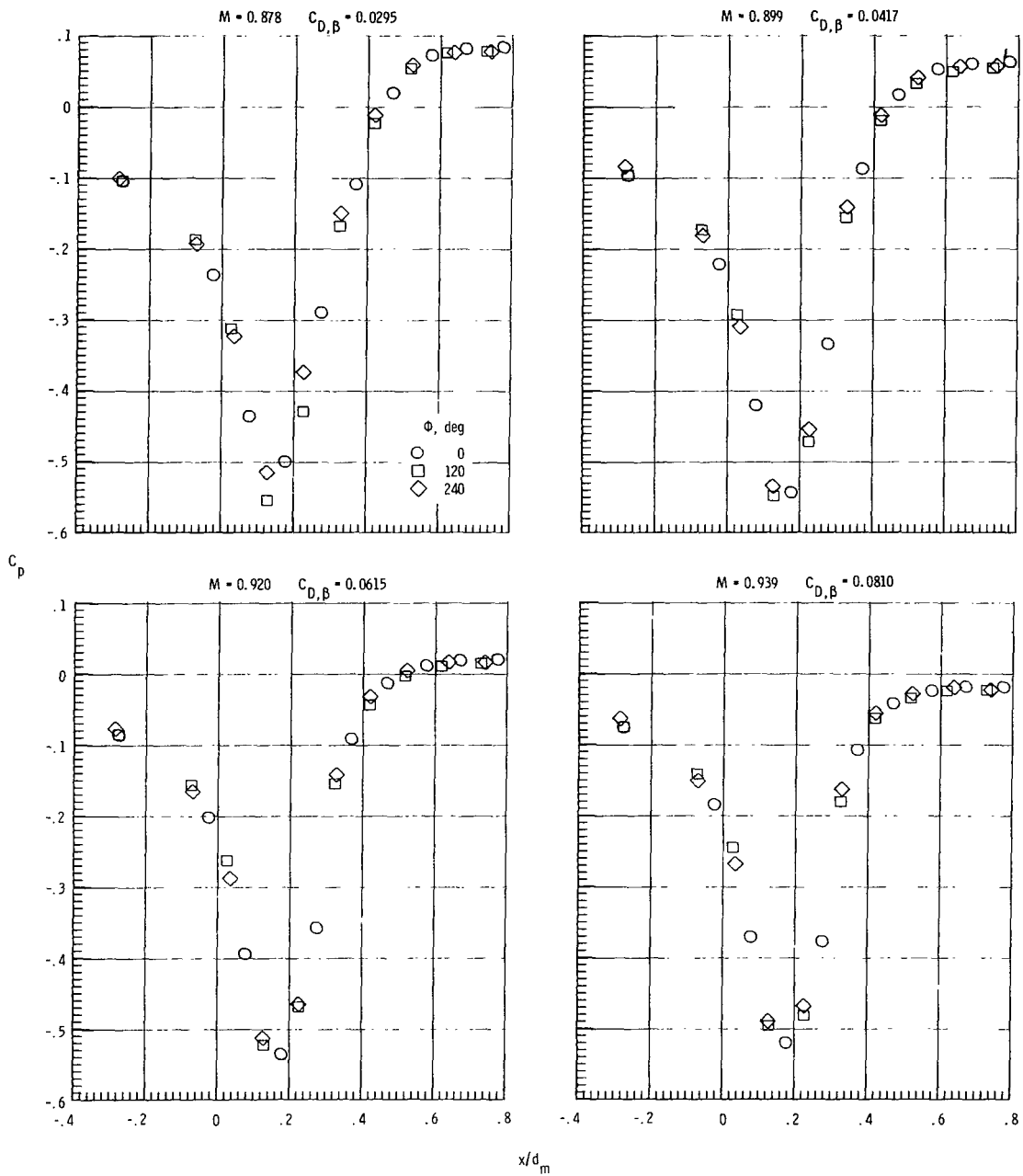
(a)  $M = 0.4$  and  $0.6$ .

Figure 5.- Boattail pressure coefficient distributions for  $L/d_m = 8.0$ ,  $l/d_m = 0.80$  circular-arc boattail at  $0^\circ$  angle of attack at various Mach numbers.



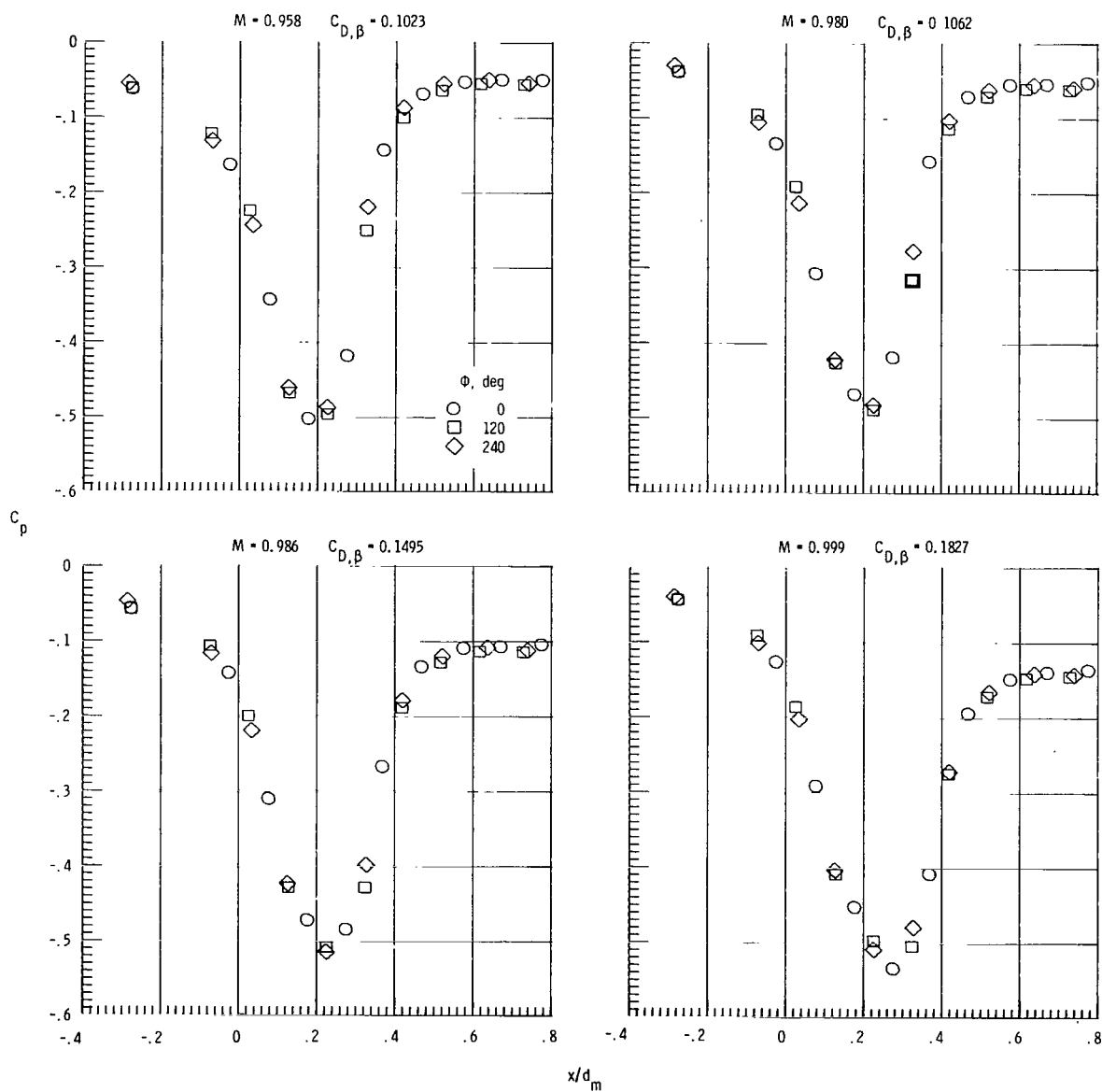
(b)  $M = 0.6$  to  $0.85$ .

Figure 5.- Continued.



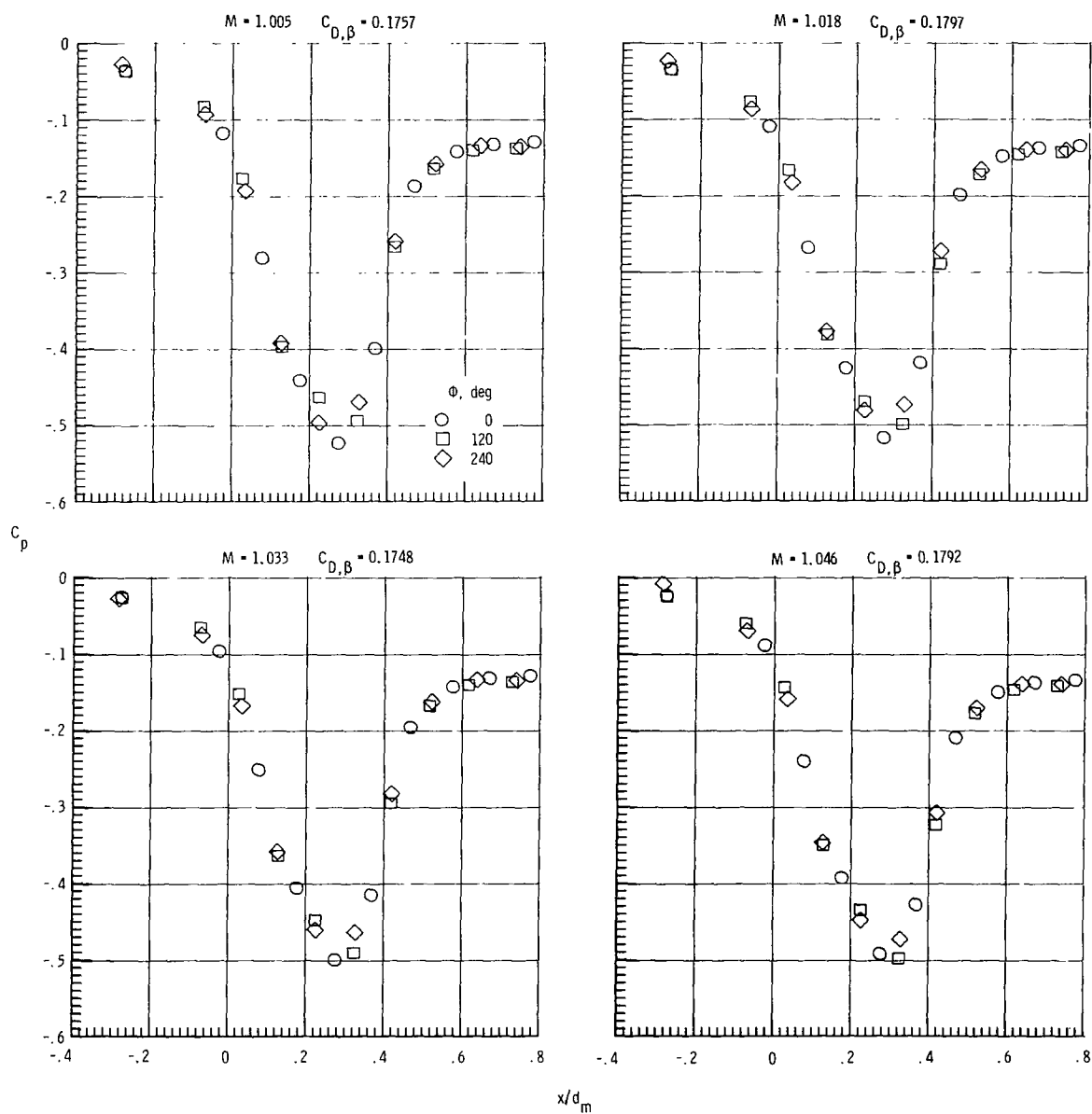
(c)  $M = 0.88$  to  $0.94$ .

Figure 5.- Continued.



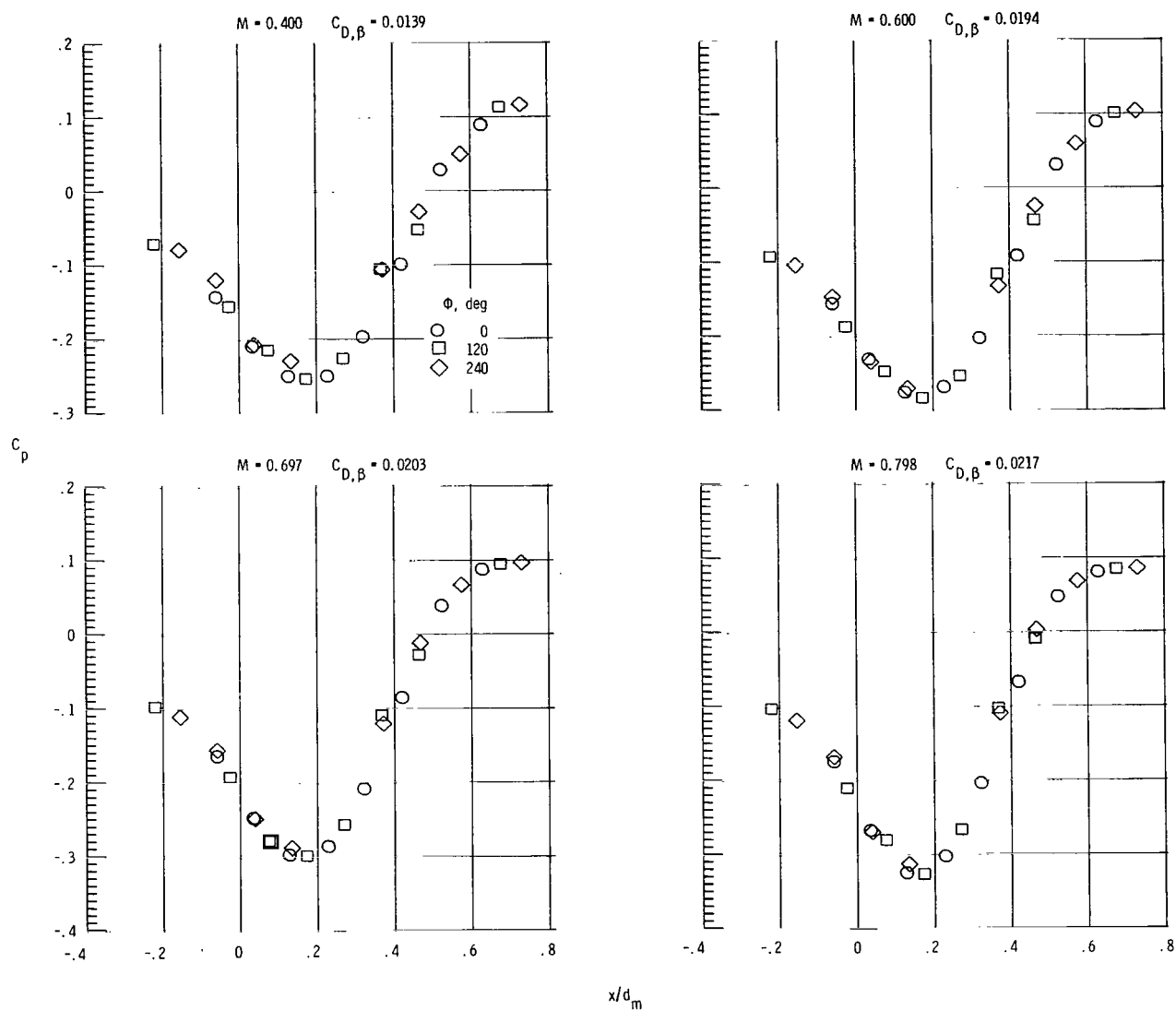
(d)  $M = 0.96$  to  $1.00$ .

Figure 5.- Continued.



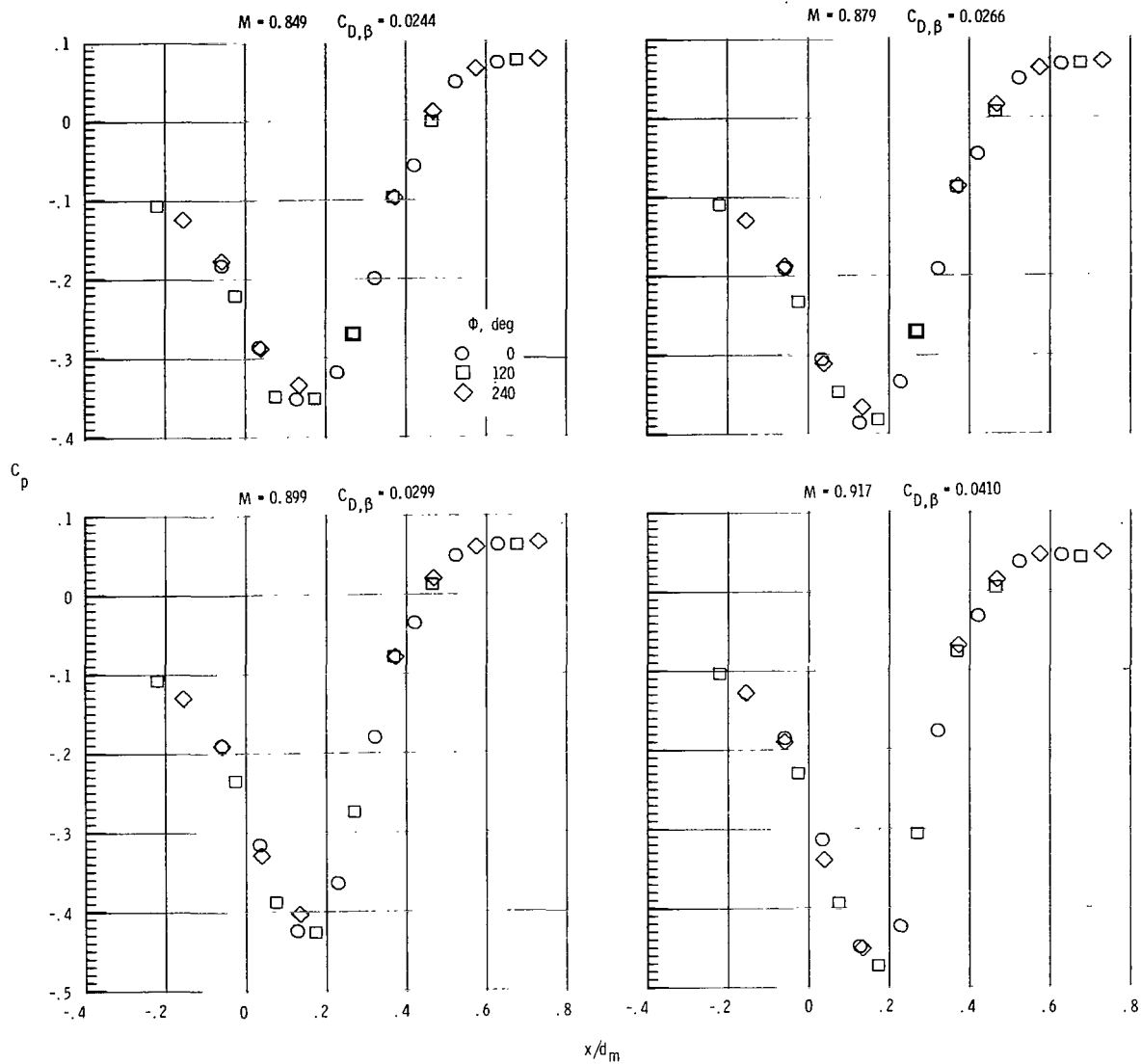
(e)  $M = 1.01$  to  $1.05$ .

Figure 5.- Concluded.



(a)  $M = 0.4$  to  $0.8$ .

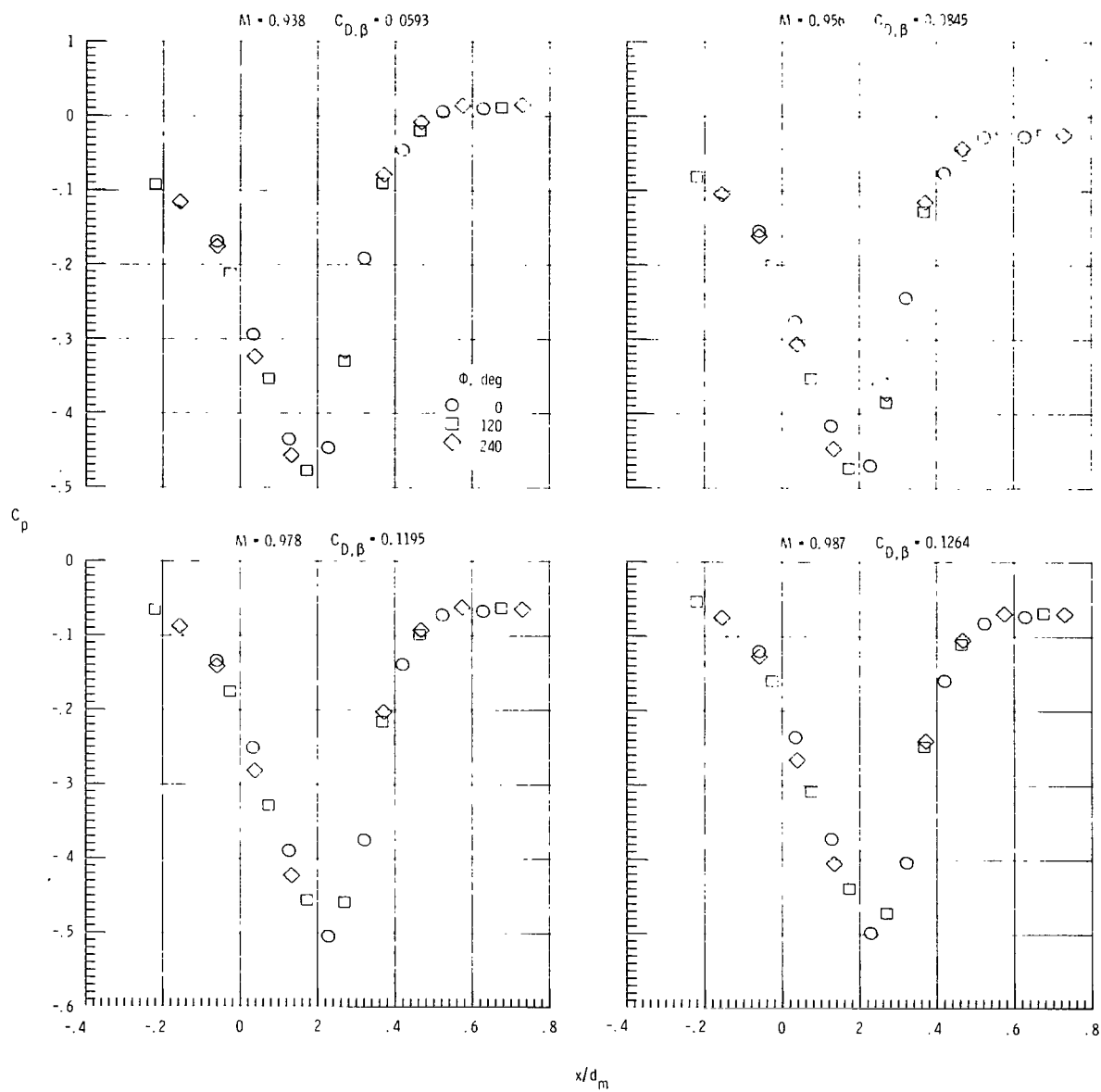
Figure 6.- Boattail pressure coefficient distributions for  $L/d_m = 16.0$ ,  $l/d_m = 0.80$  circular-arc boattail at  $0^\circ$  angle of attack at various Mach numbers.



(b)  $M = 0.85$  to  $0.92$ .

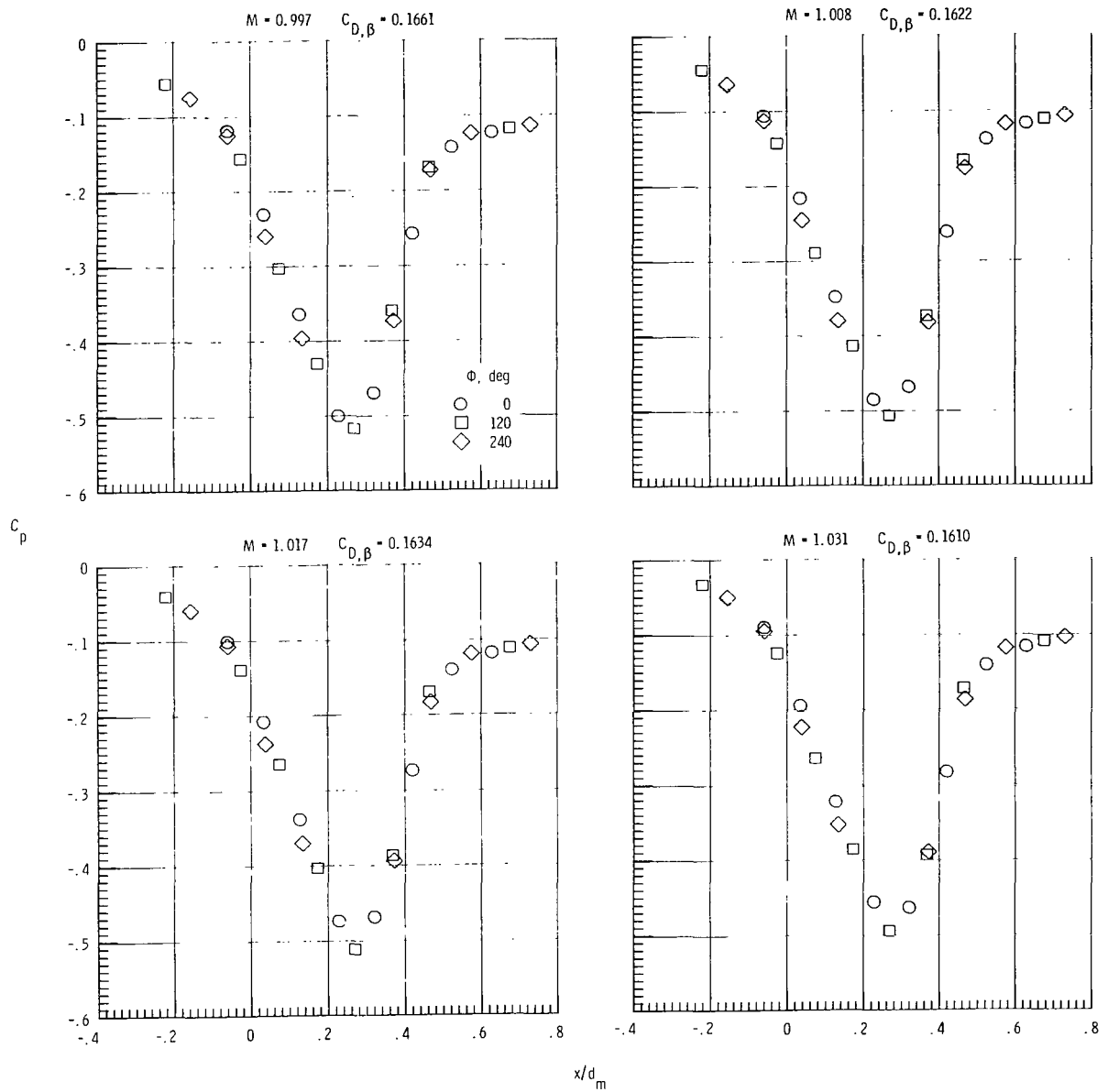
Figure 6.- Continued.





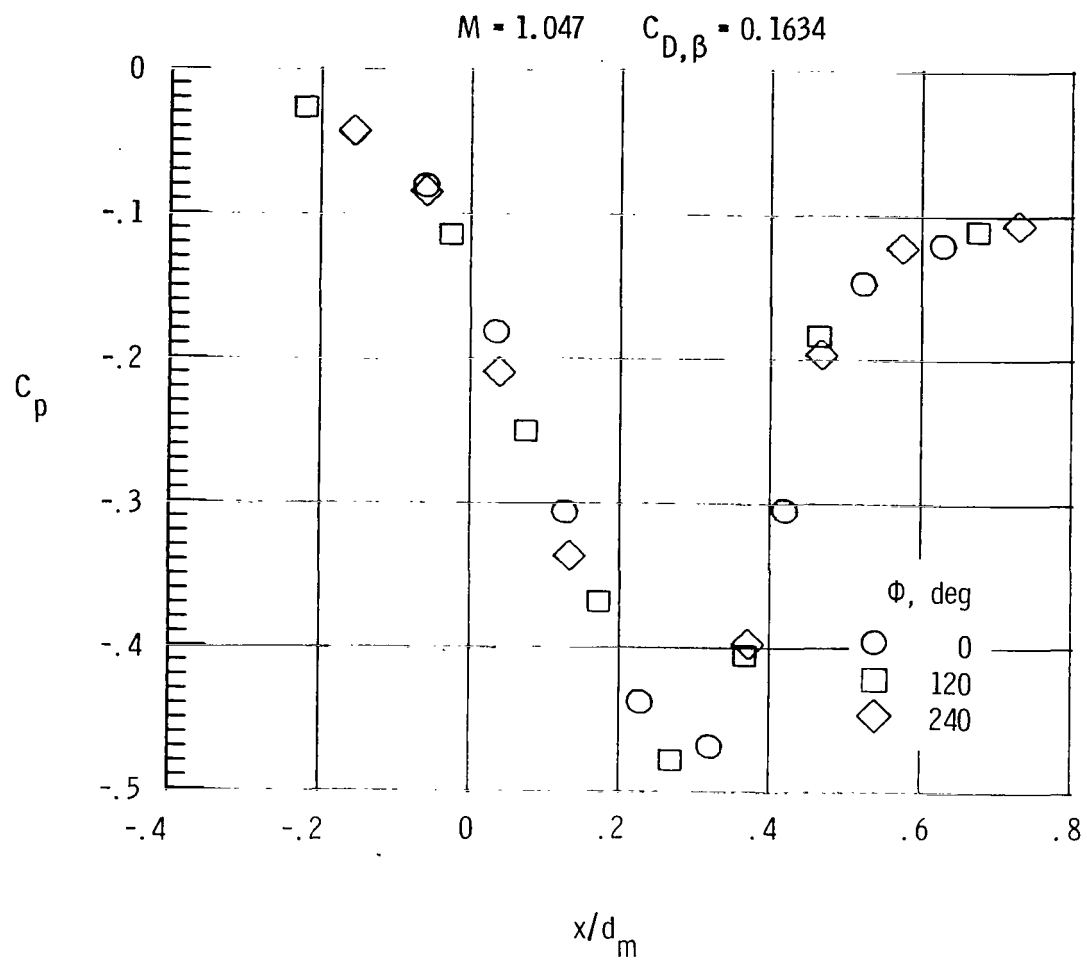
(c)  $M = 0.94$  to  $0.99$ .

Figure 6.- Continued.



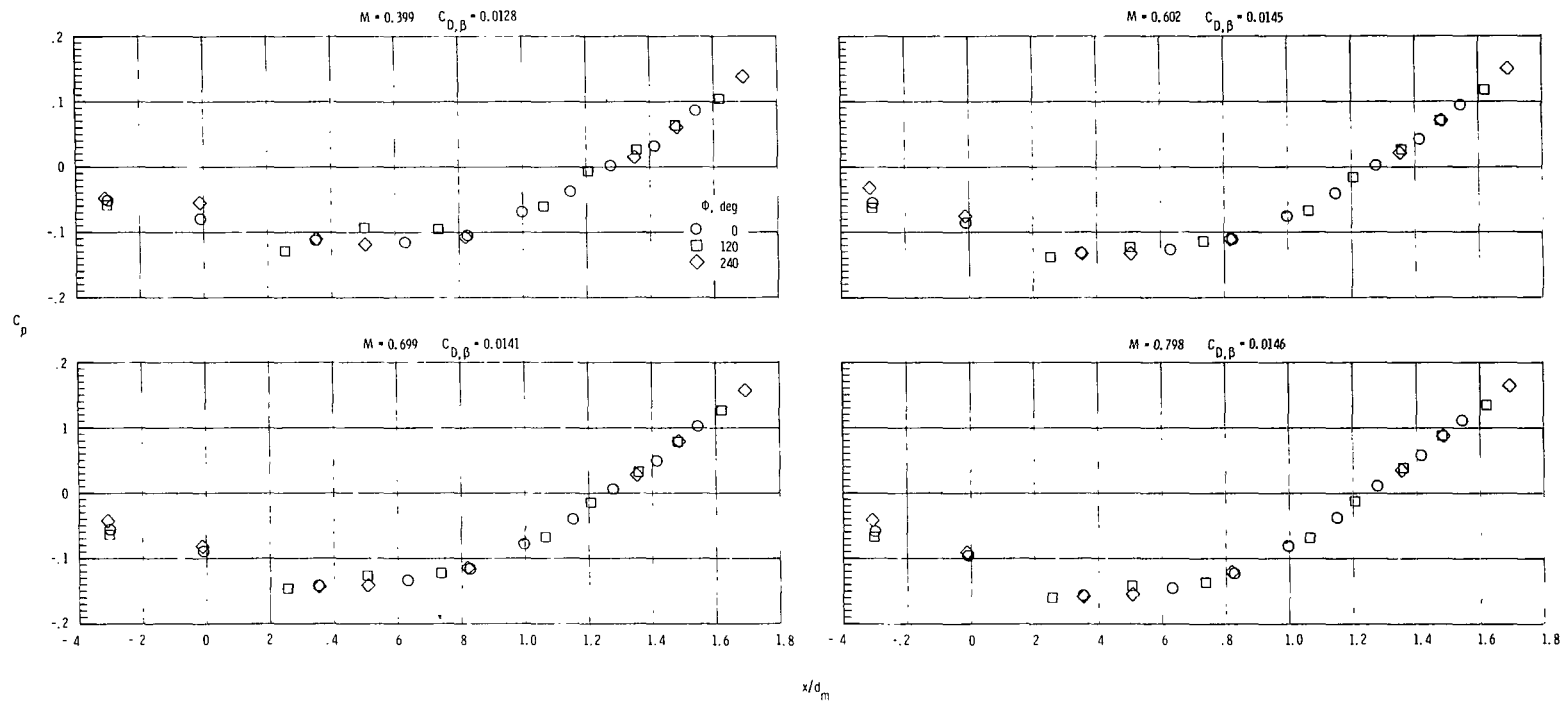
(d)  $M = 1.00$  to  $1.03$ .

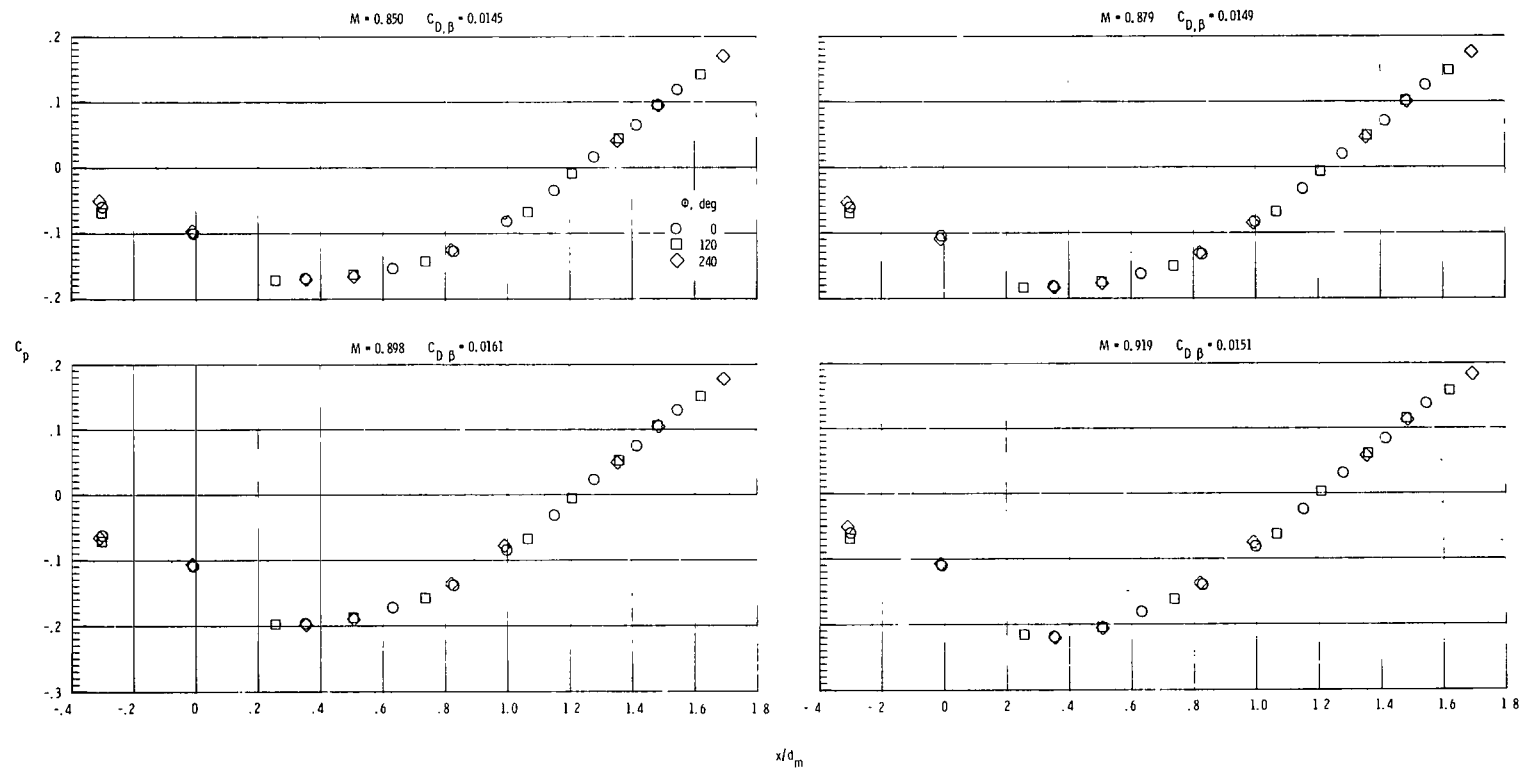
Figure 6.- Continued.



(e)  $M = 1.05$ .

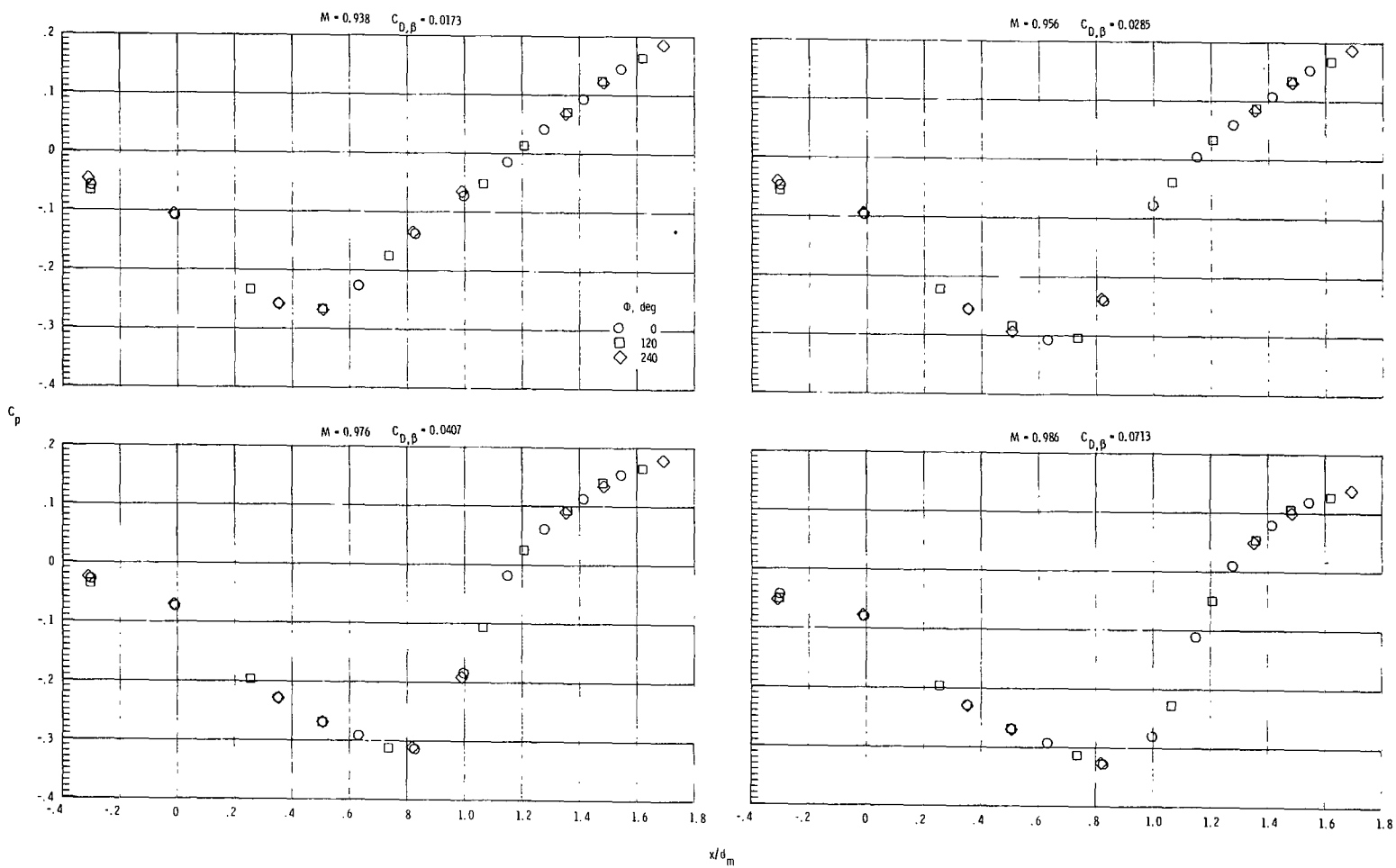
Figure 6.- Concluded.

(a)  $M = 0.4$  to  $0.8$ .Figure 7.- Boattail pressure coefficient distributions for  $L/d_m = 8.0$ ,  $l/d_m = 1.77$  circular-arc boattail at  $0^\circ$  angle of attack at various Mach numbers.



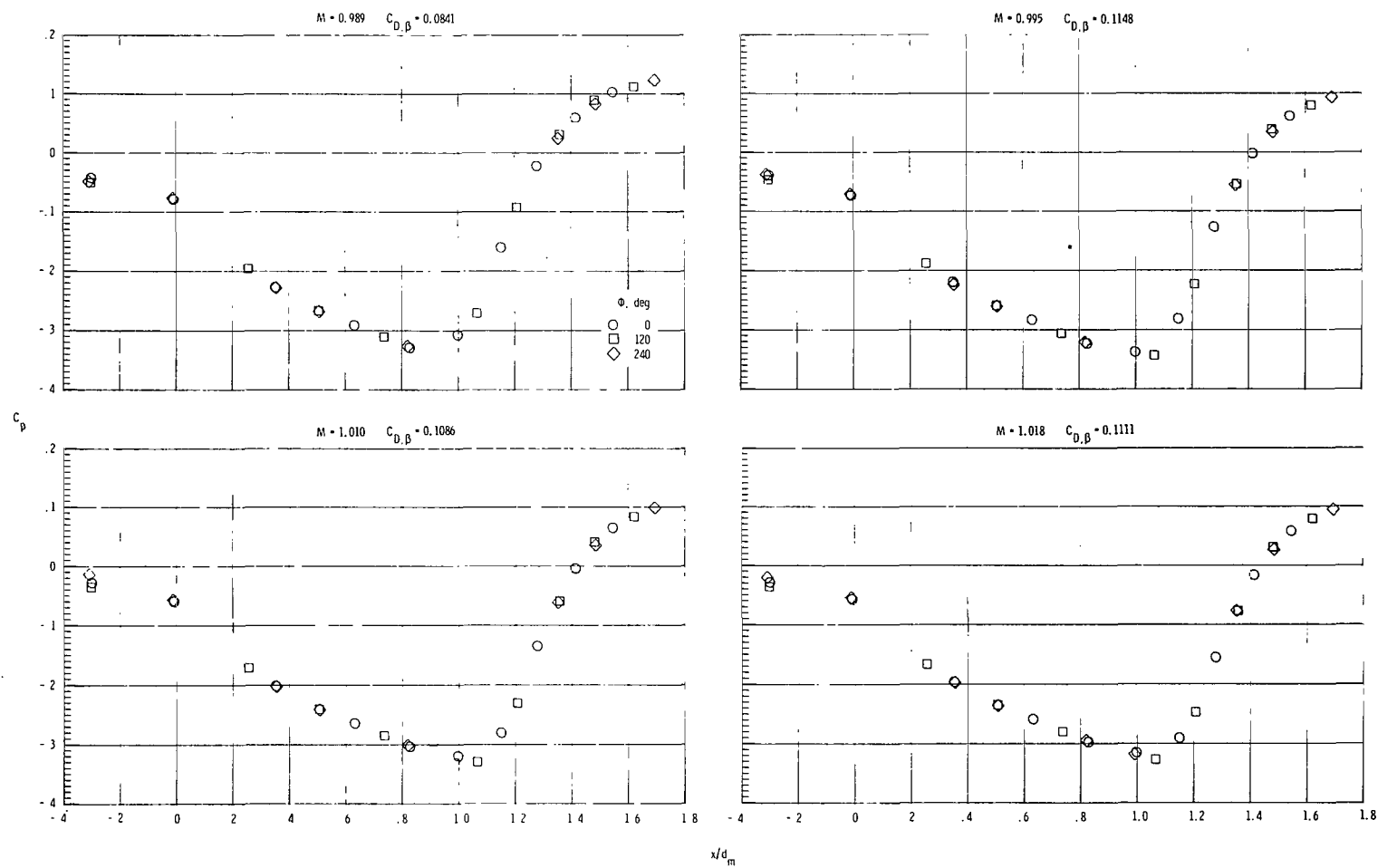
(b)  $M = 0.85$  to  $0.92$ .

Figure 7.- Continued.



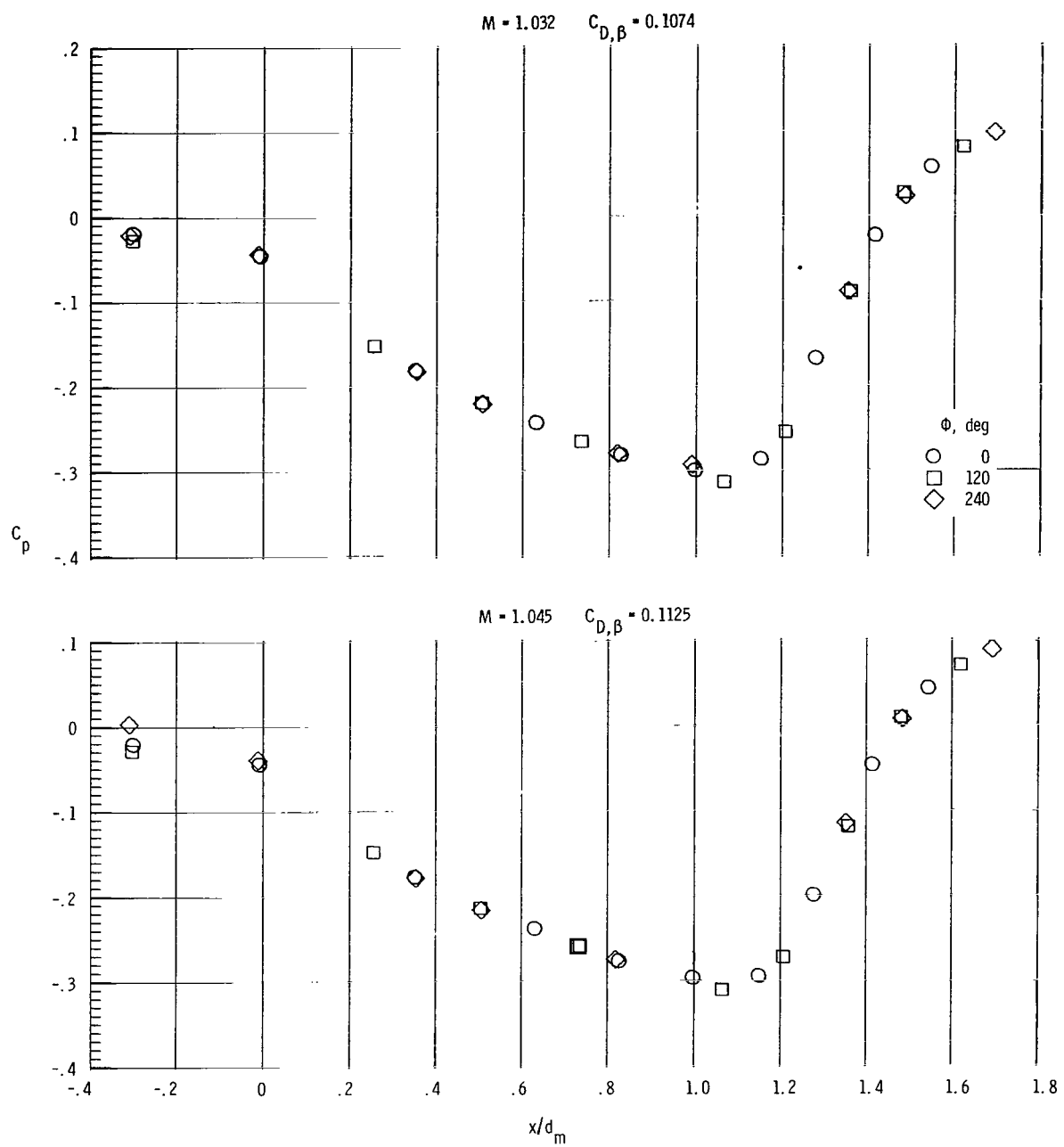
(c)  $M = 0.94$  to  $0.99$ .

Figure 7.- Continued.



(d)  $M = 0.99$  to  $1.02$ .

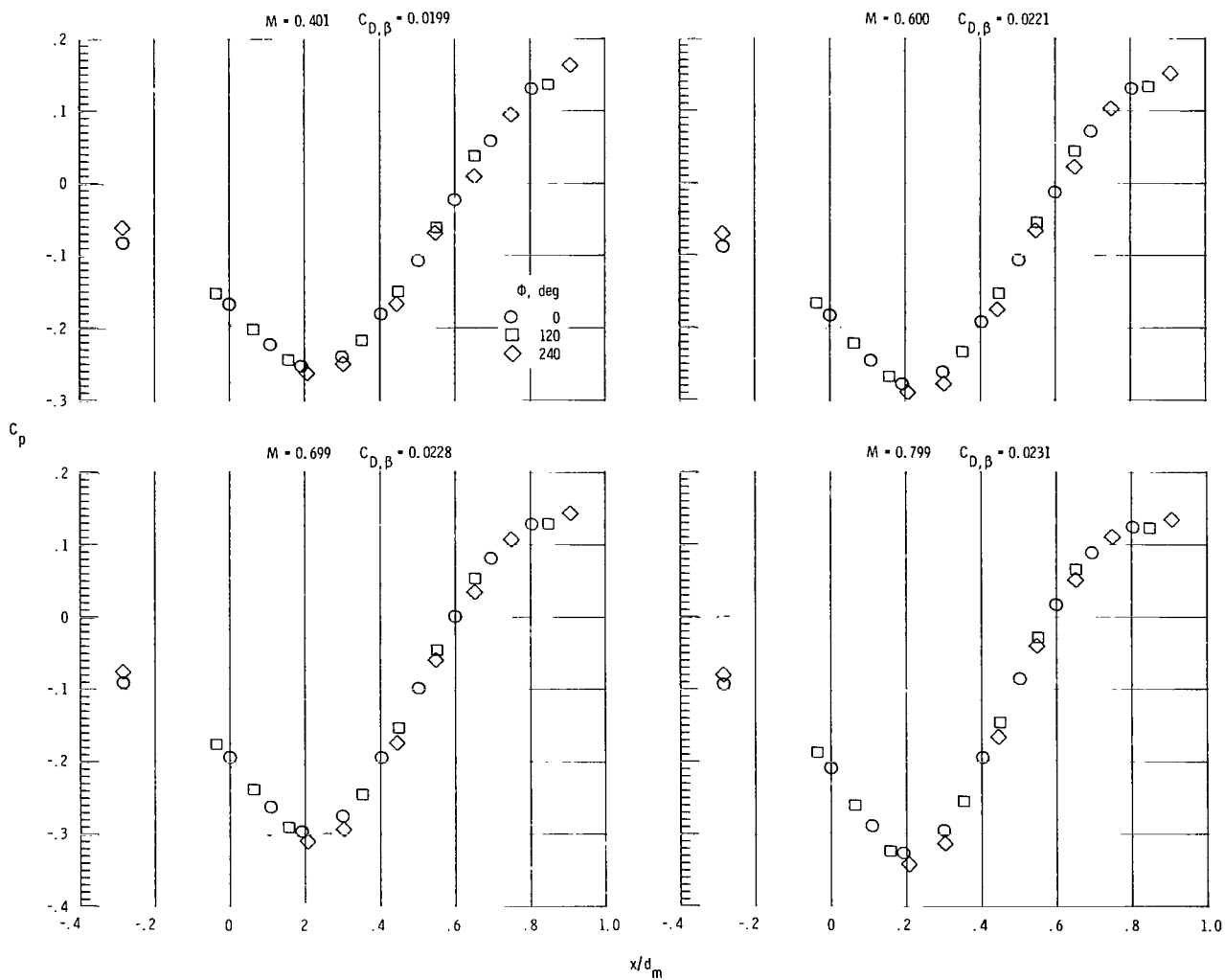
Figure 7.- Continued.



(e)  $M = 1.03$  and  $1.05$ .

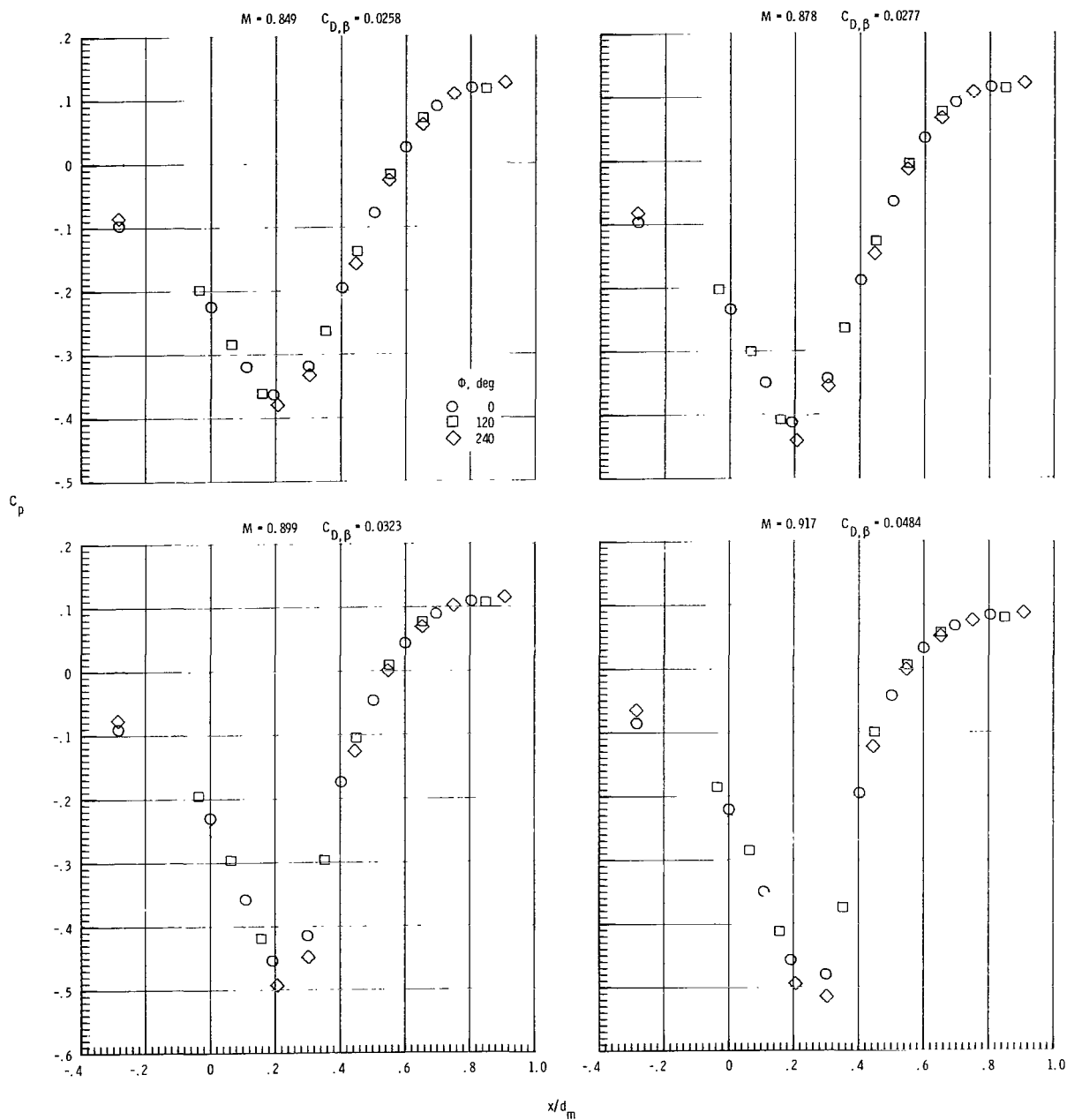
Figure 7.- Concluded.





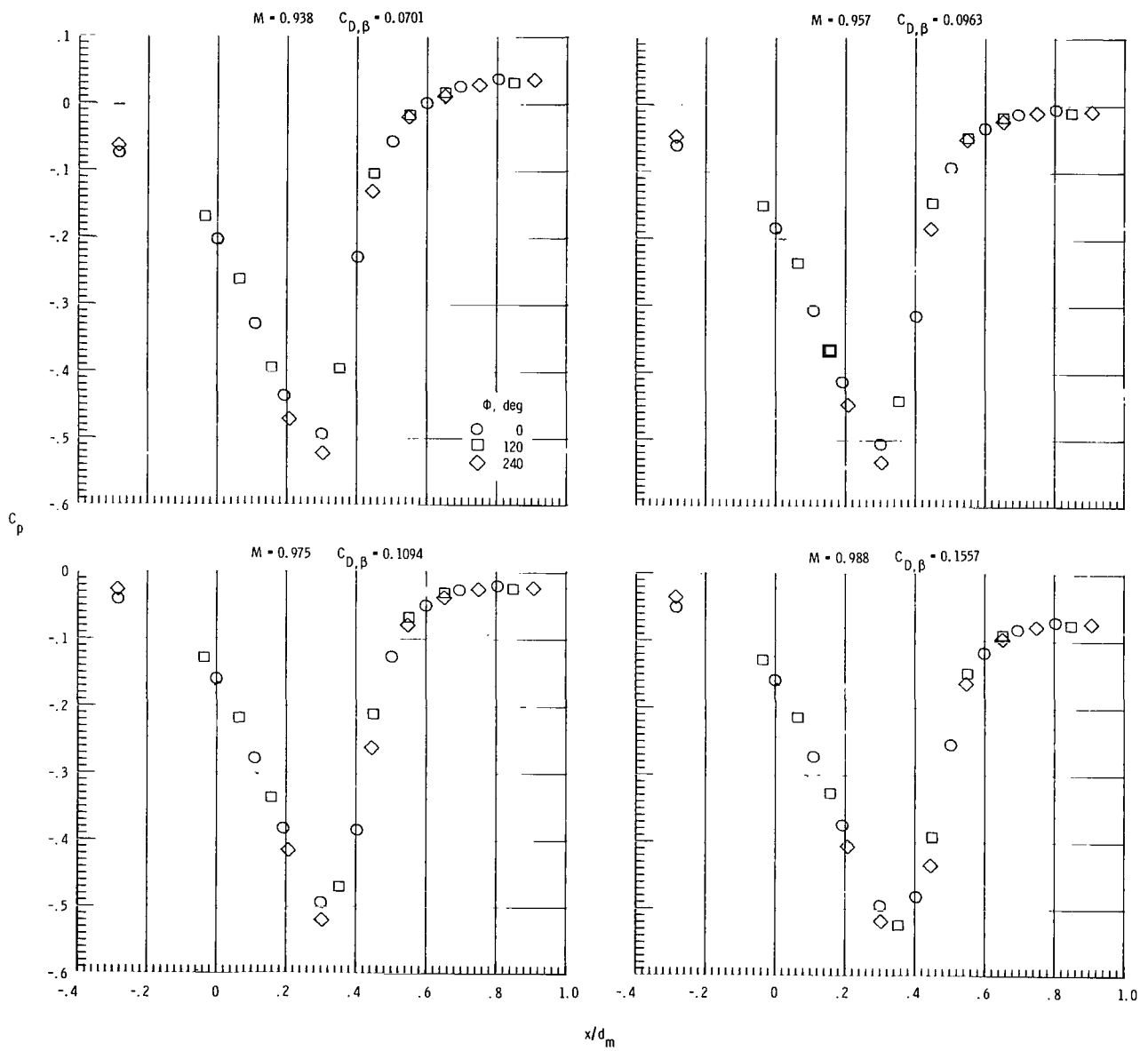
(a)  $M = 0.4$  to  $0.8$ .

Figure 8.- Boattail pressure coefficient distributions for  $L/d_m = 8.0$ ,  $l/d_m = 0.96$  circular-arc-conic boattail at  $0^\circ$  angle of attack at various Mach numbers.



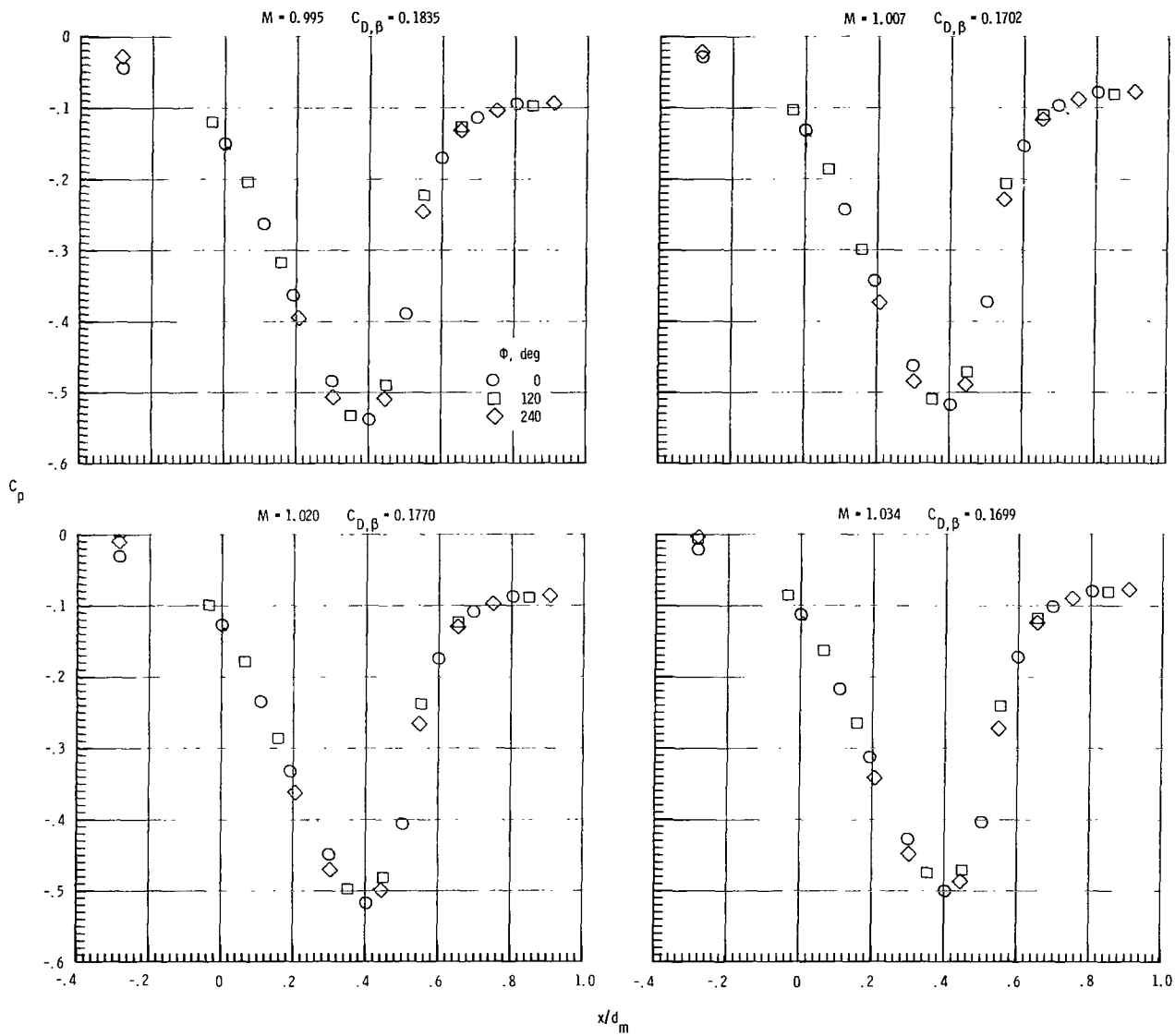
(b)  $M = 0.85$  to  $0.92$ .

Figure 8.- Continued.



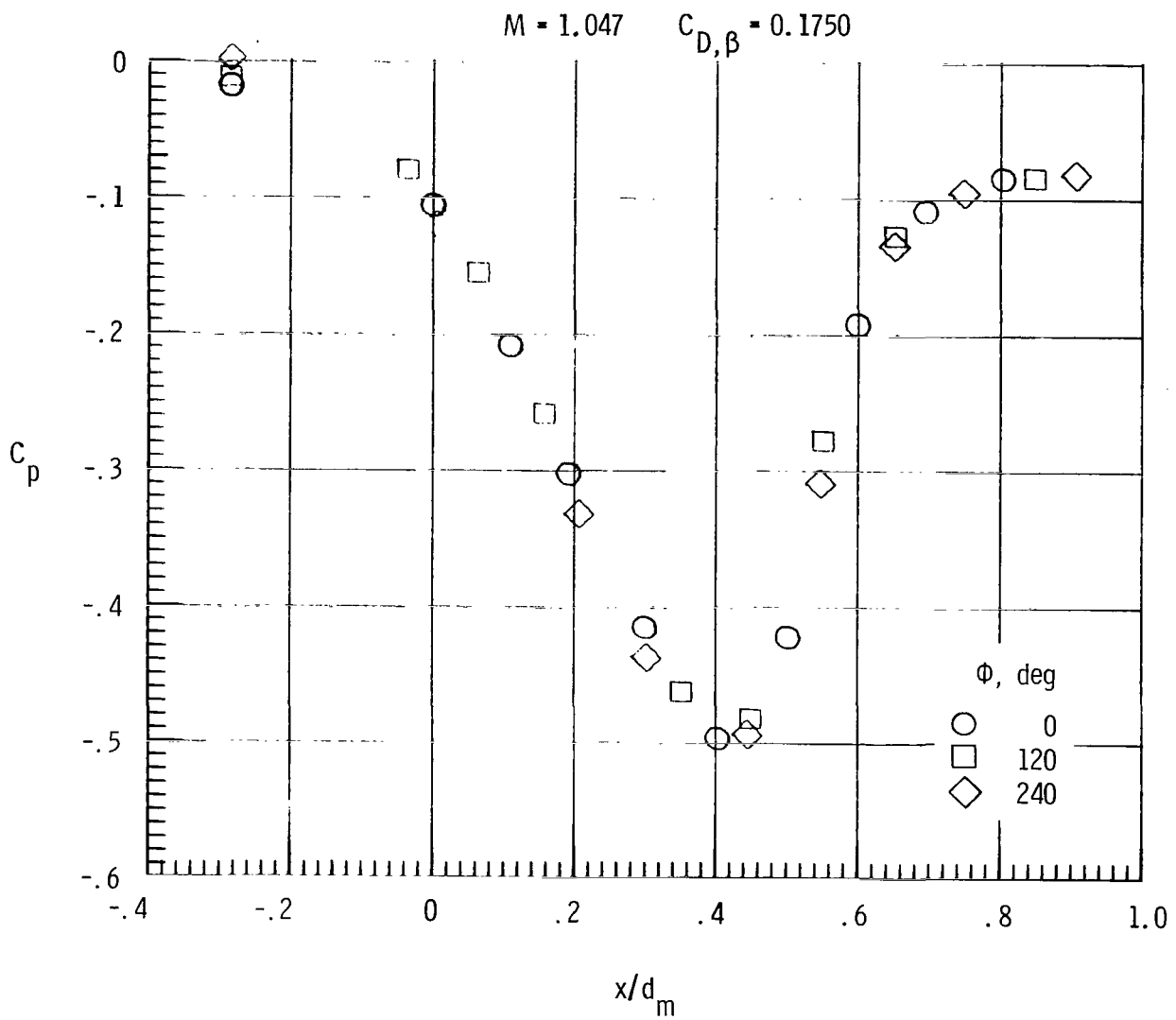
(c)  $M = 0.94$  to  $0.99$ .

Figure 8.- Continued.



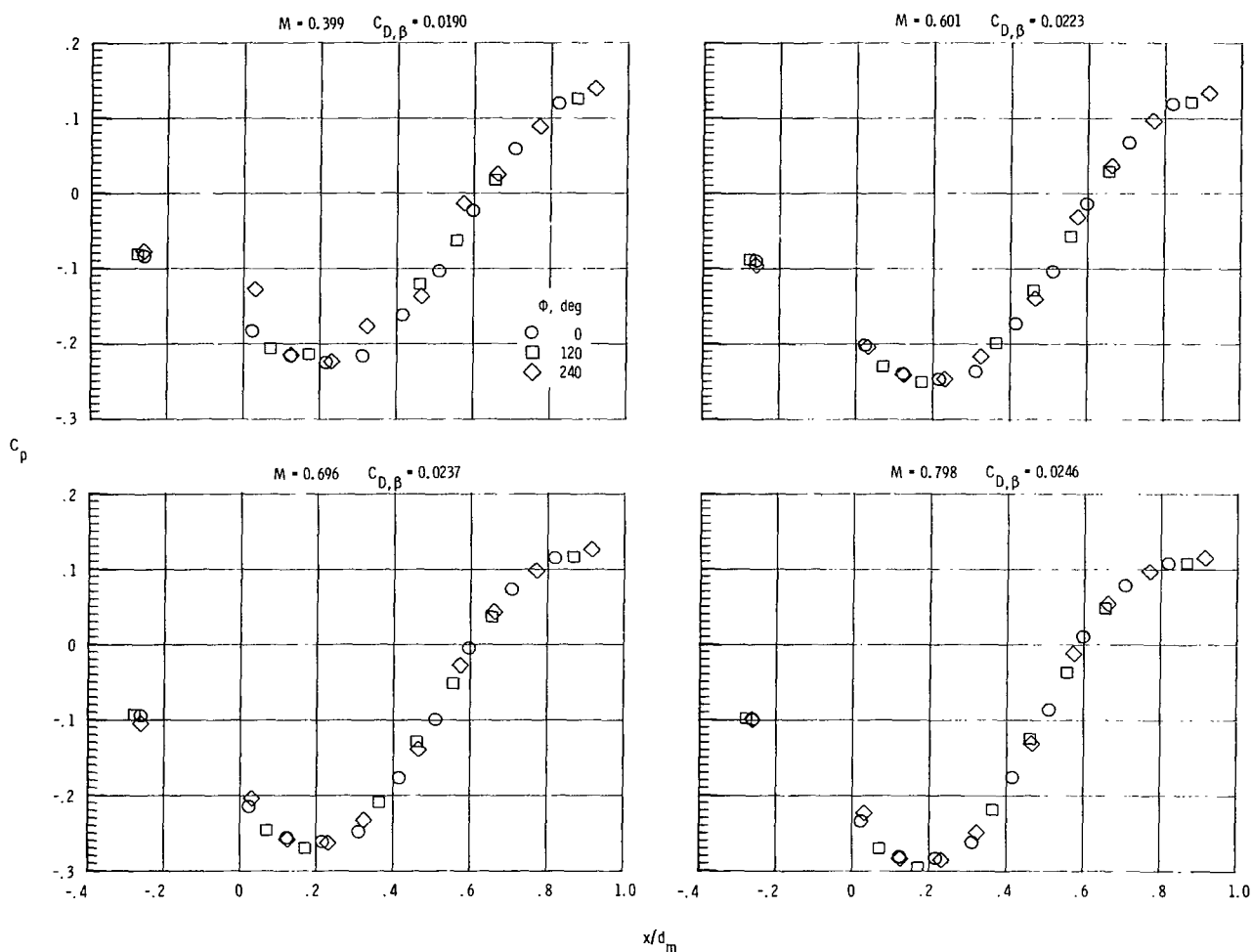
(d)  $M = 1.00$  to  $1.03$ .

Figure 8.- Continued.



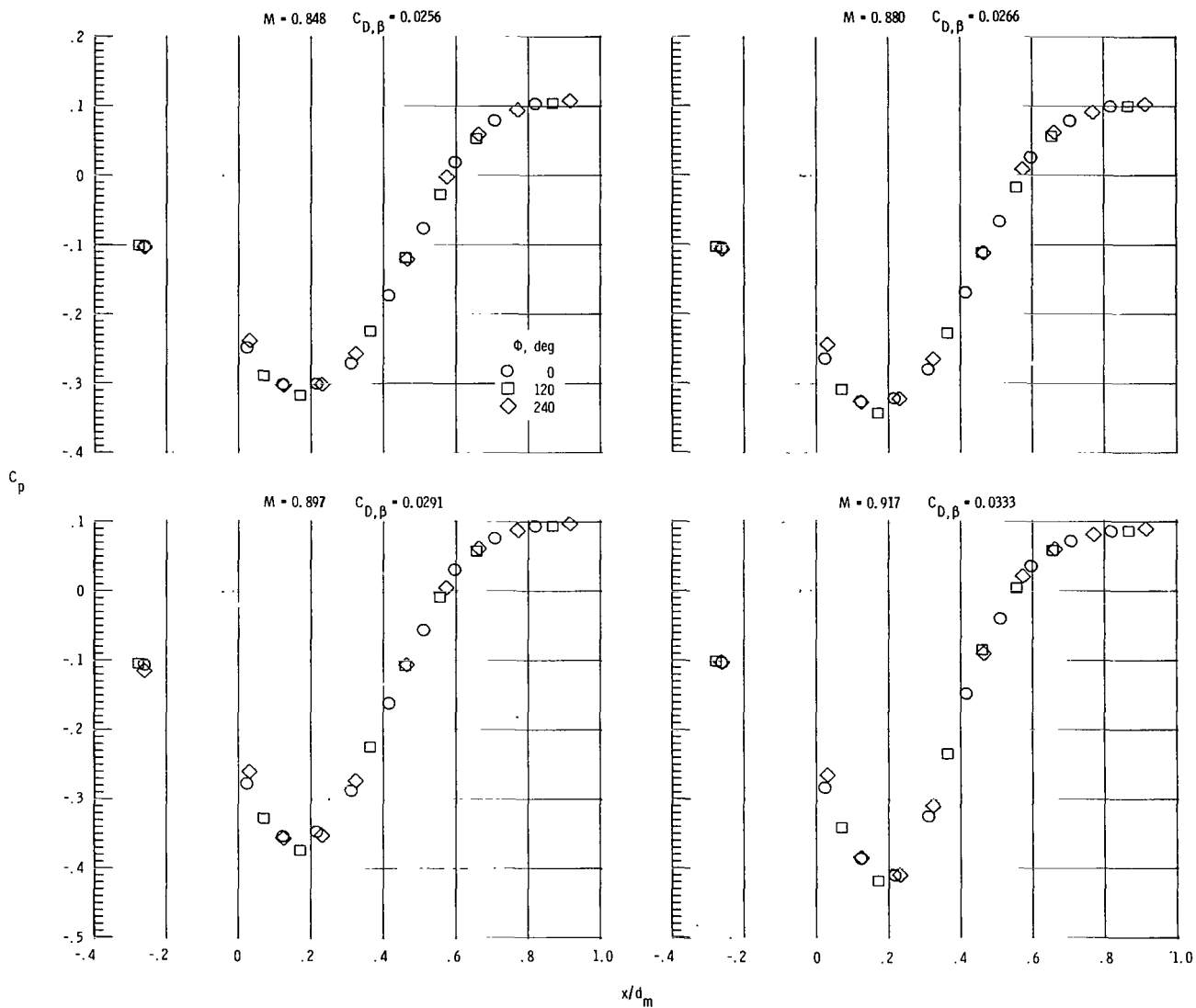
(e)  $M = 1.05$ .

Figure 8.- Concluded.



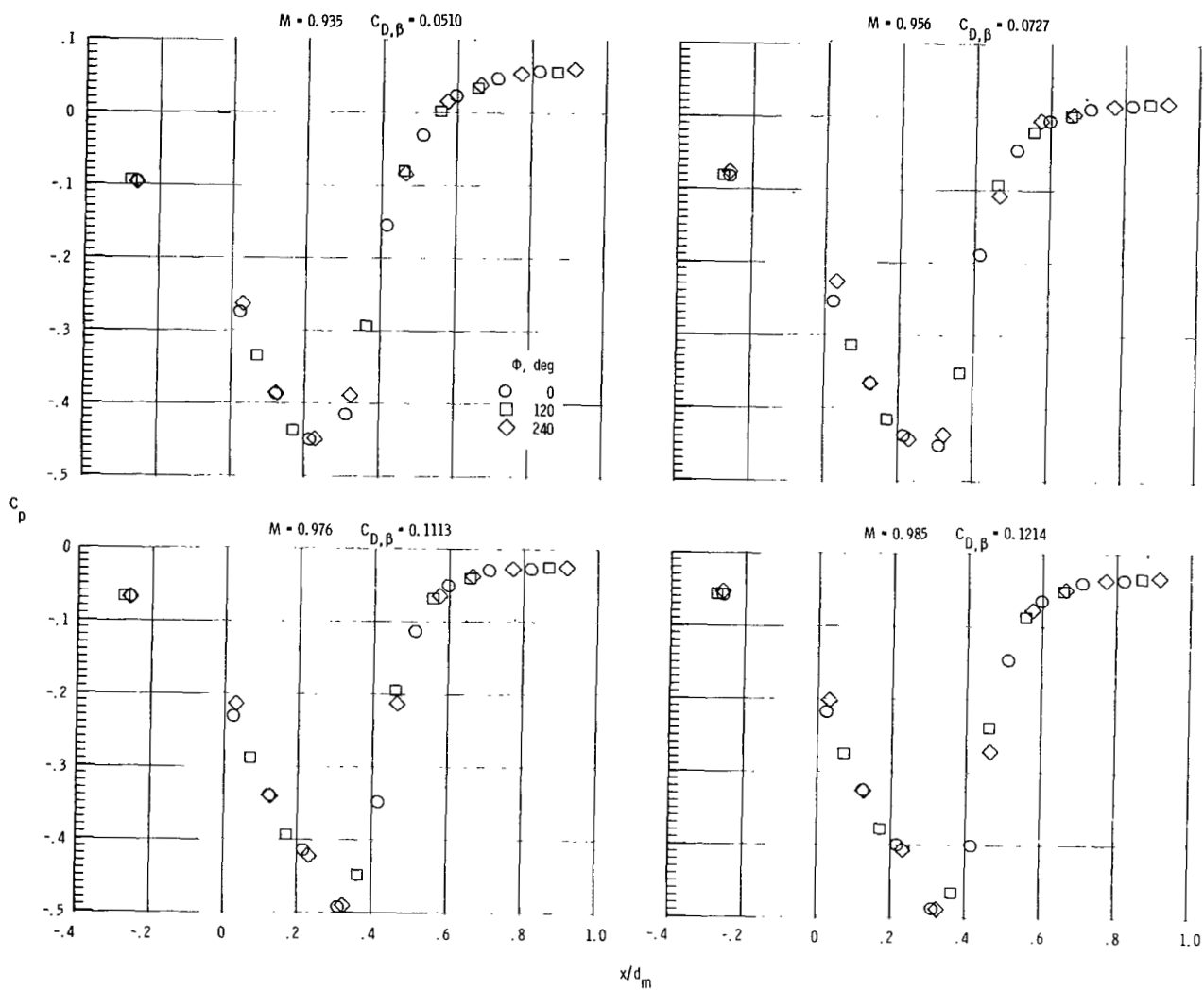
(a)  $M = 0.4$  to  $0.8$ .

Figure 9.- Boattail pressure coefficient distributions for  $L/d_m = 16.0$ ,  $l/d_m = 0.96$  circular-arc—conic boattail at  $0^\circ$  angle of attack at various Mach numbers.



(b)  $M = 0.85$  to  $0.92$ .

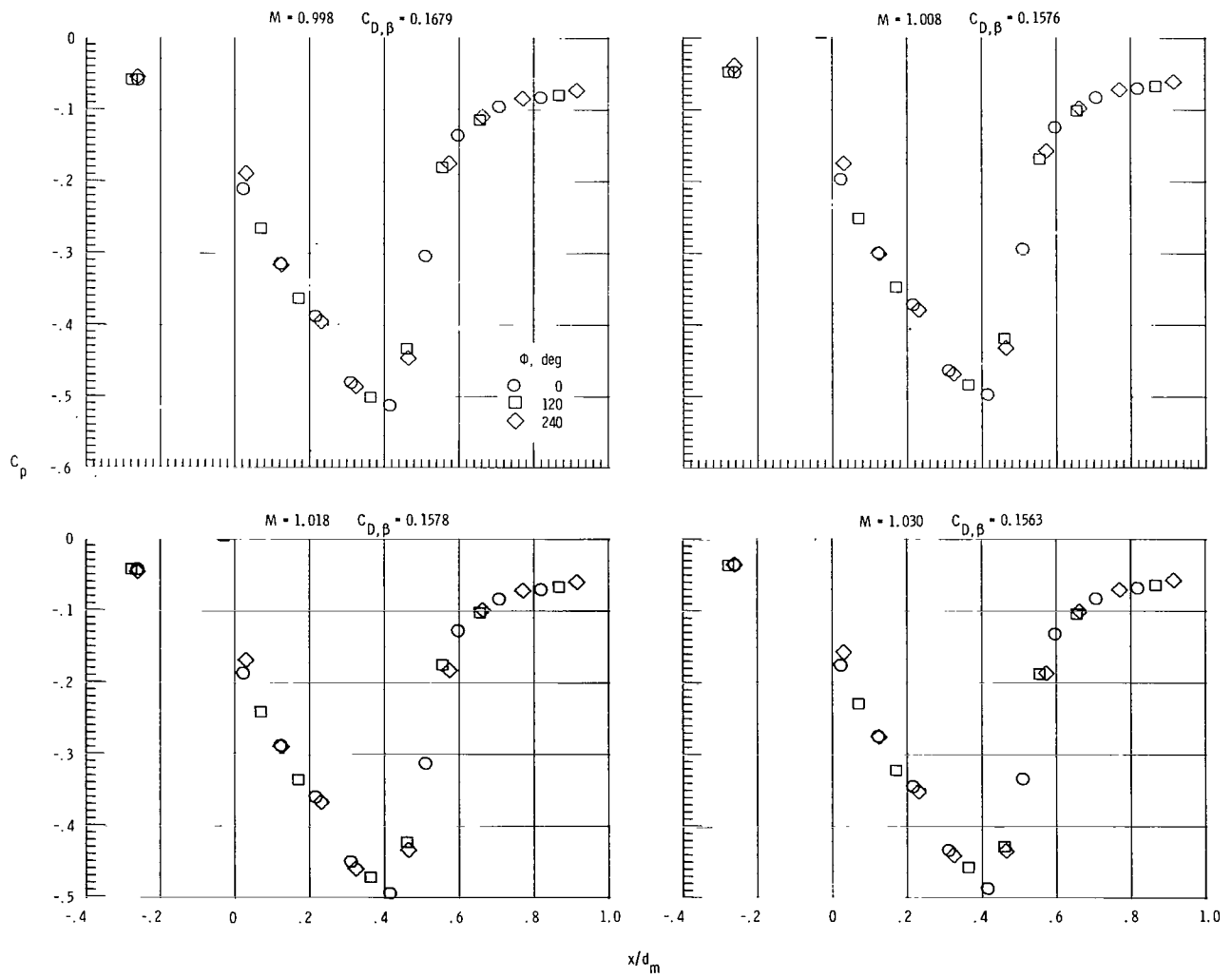
Figure 9.- Continued.



(c)  $M = 0.94$  to  $0.99$ .

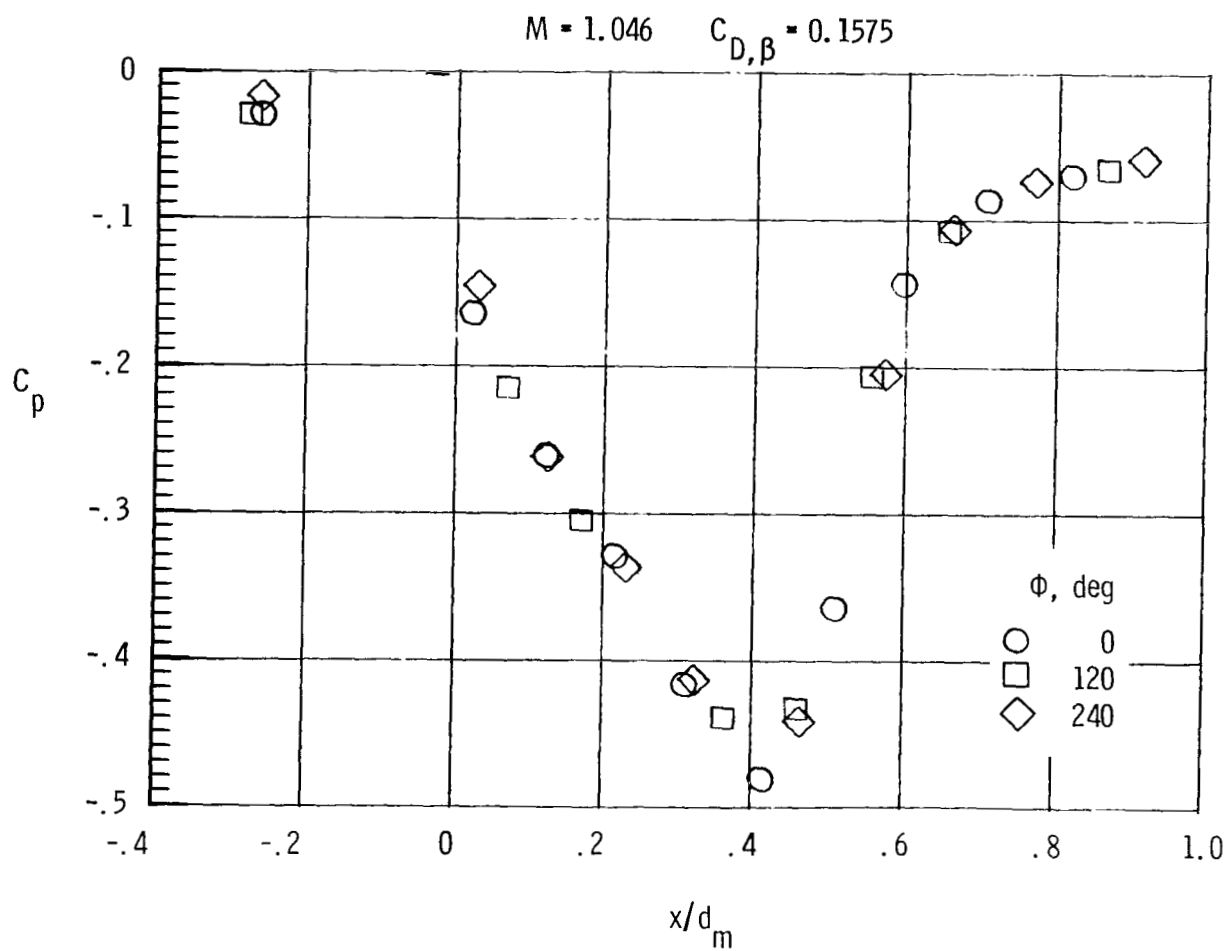
Figure 9.- Continued.





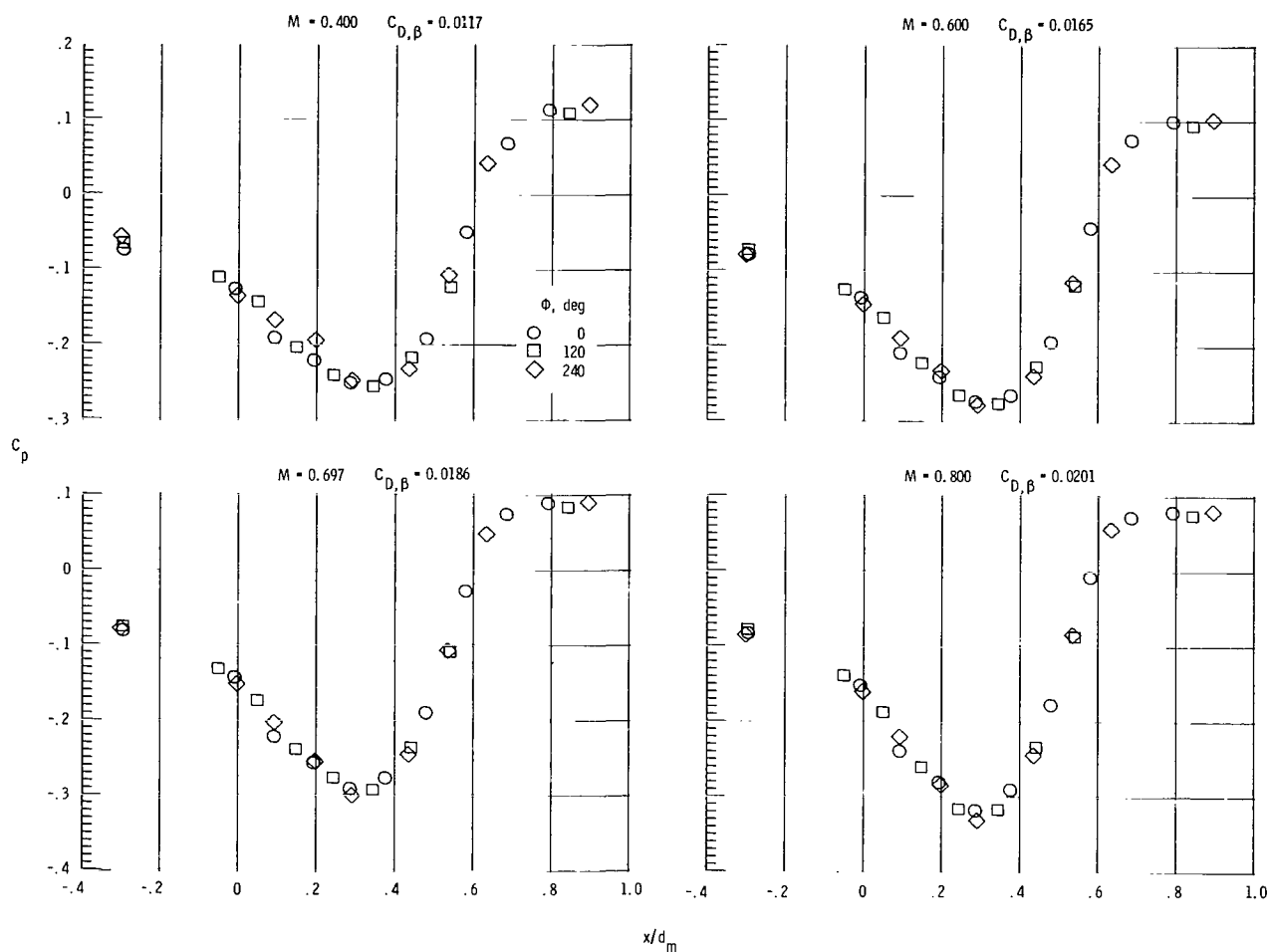
(d)  $M = 1.00$  to  $1.03$ .

Figure 9.- Continued.



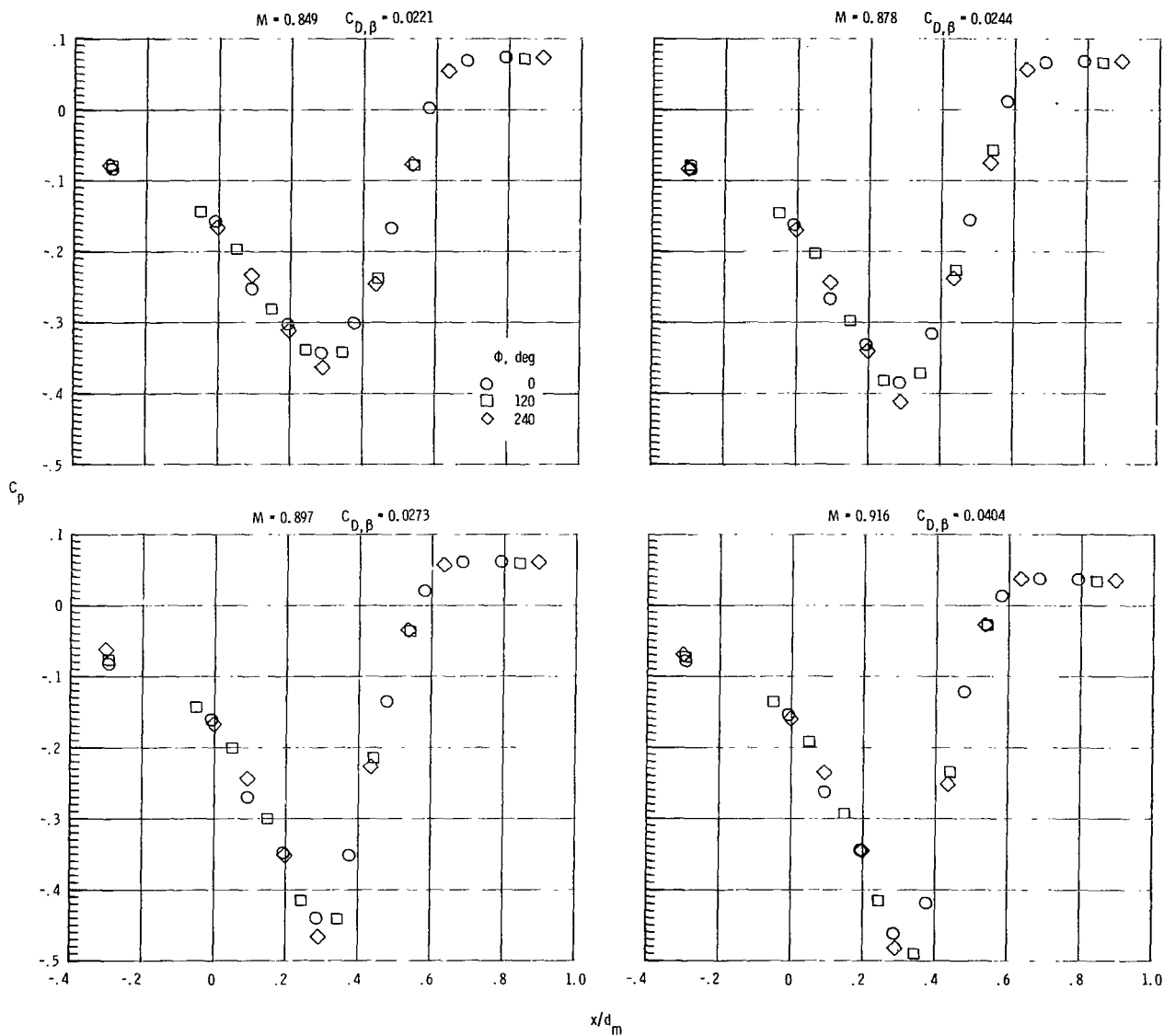
(e)  $M = 1.05$ .

Figure 9.- Concluded.



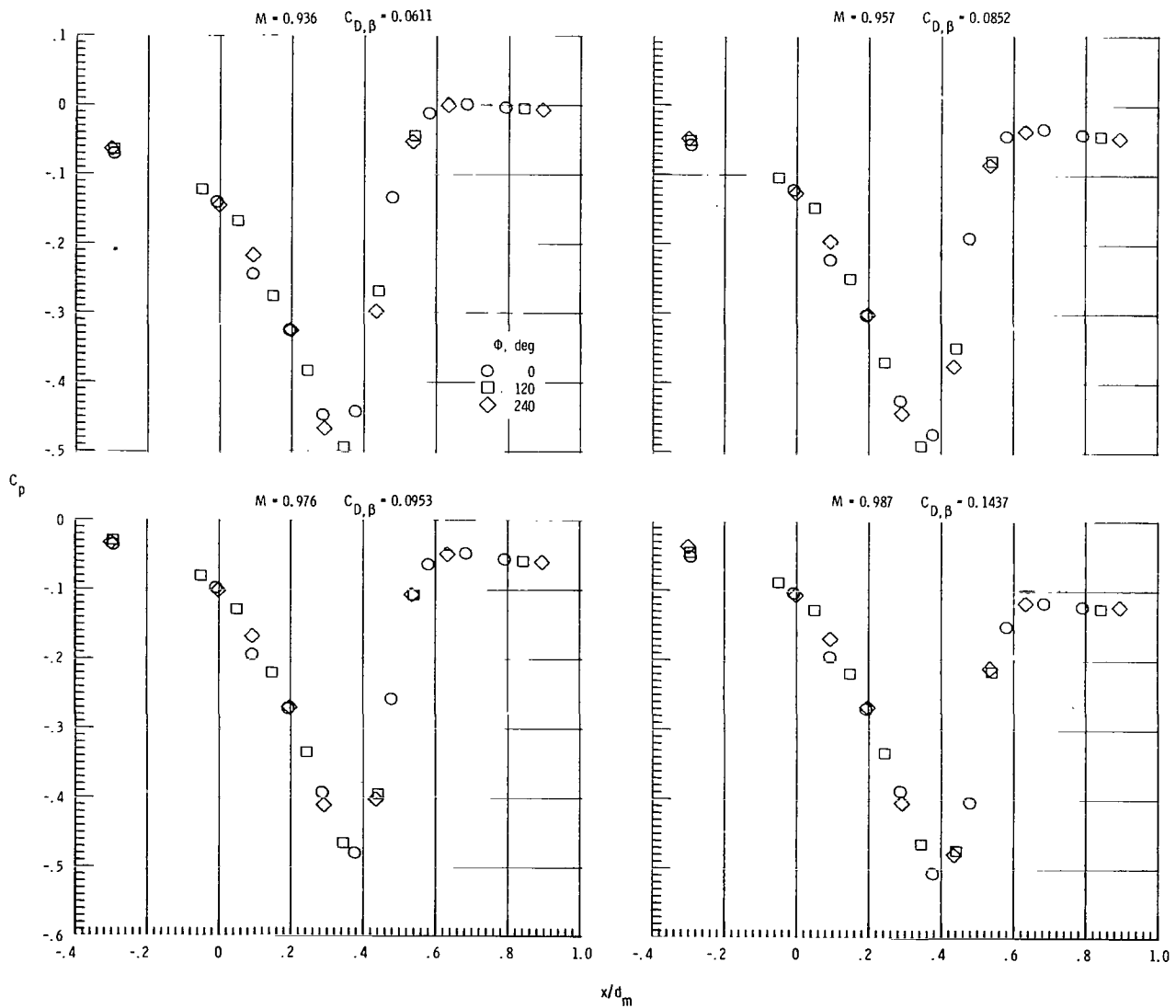
(a)  $M = 0.4$  to  $0.8$ .

Figure 10.- Boattail pressure coefficient distributions for  $L/d_m = 8.0$ ,  $l/d_m = 0.95$  contoured boattail at  $0^\circ$  angle of attack at various Mach numbers.



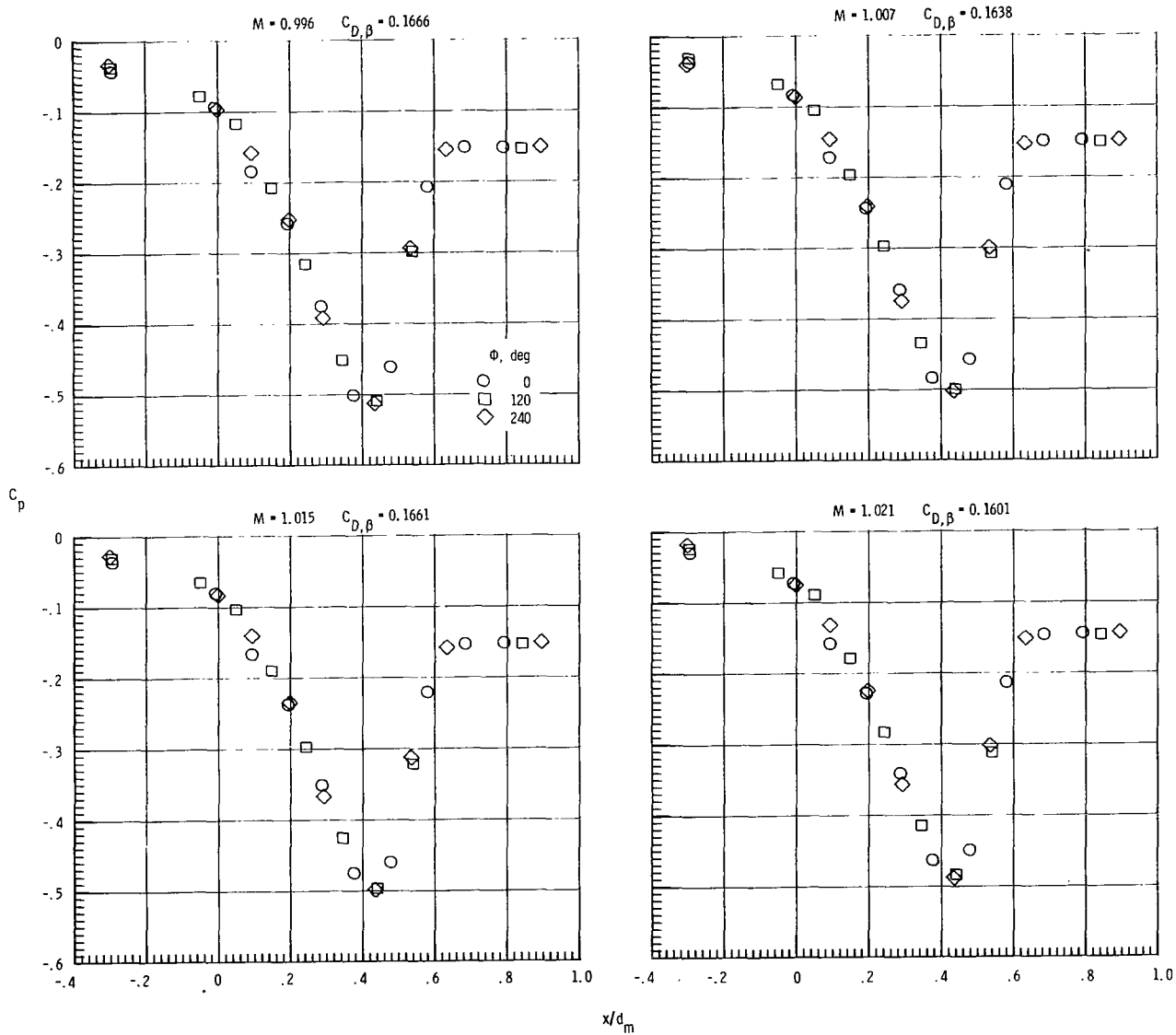
(b)  $M = 0.85$  to  $0.92$ .

Figure 10.- Continued.



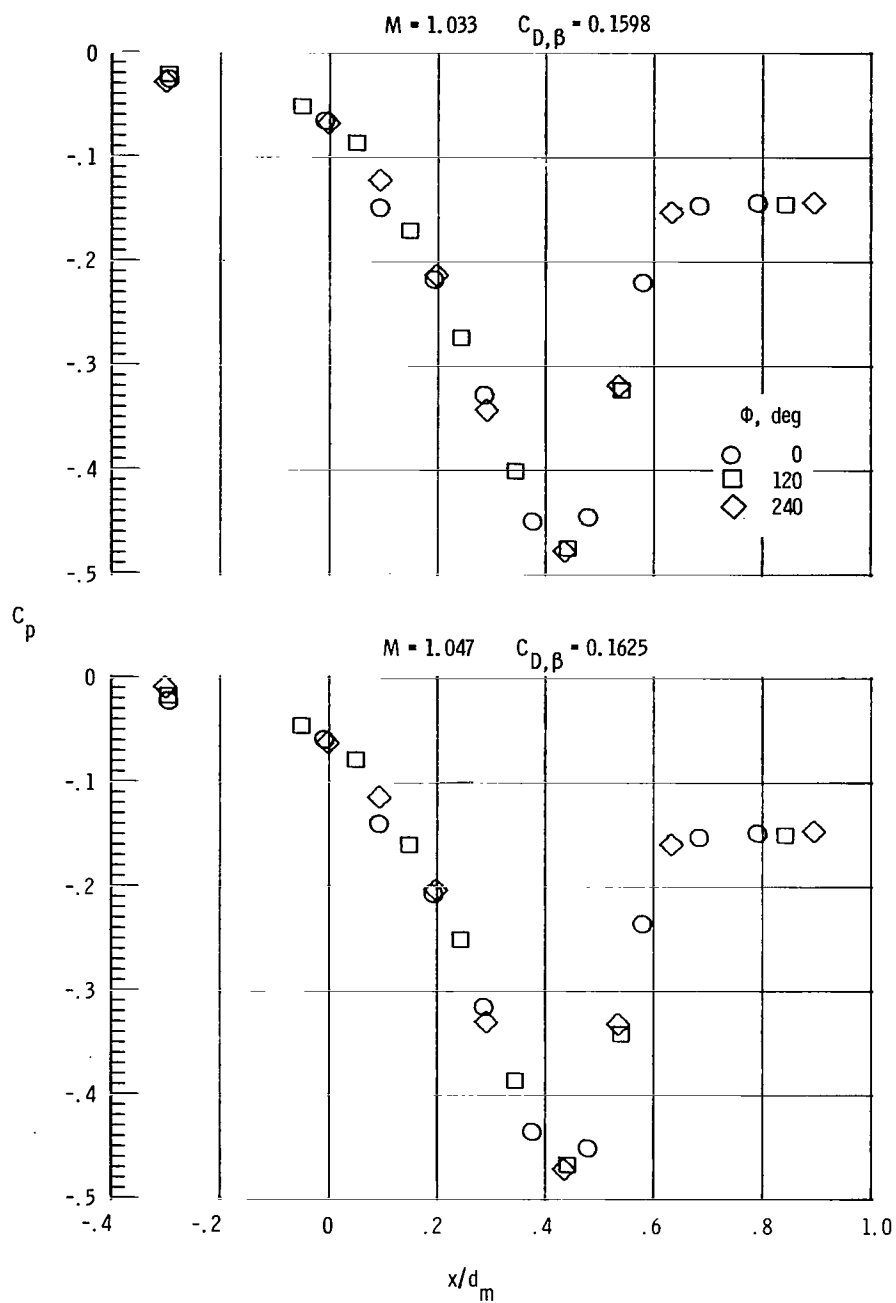
(c)  $M = 0.94$  to  $0.99$ .

Figure 10.- Continued.



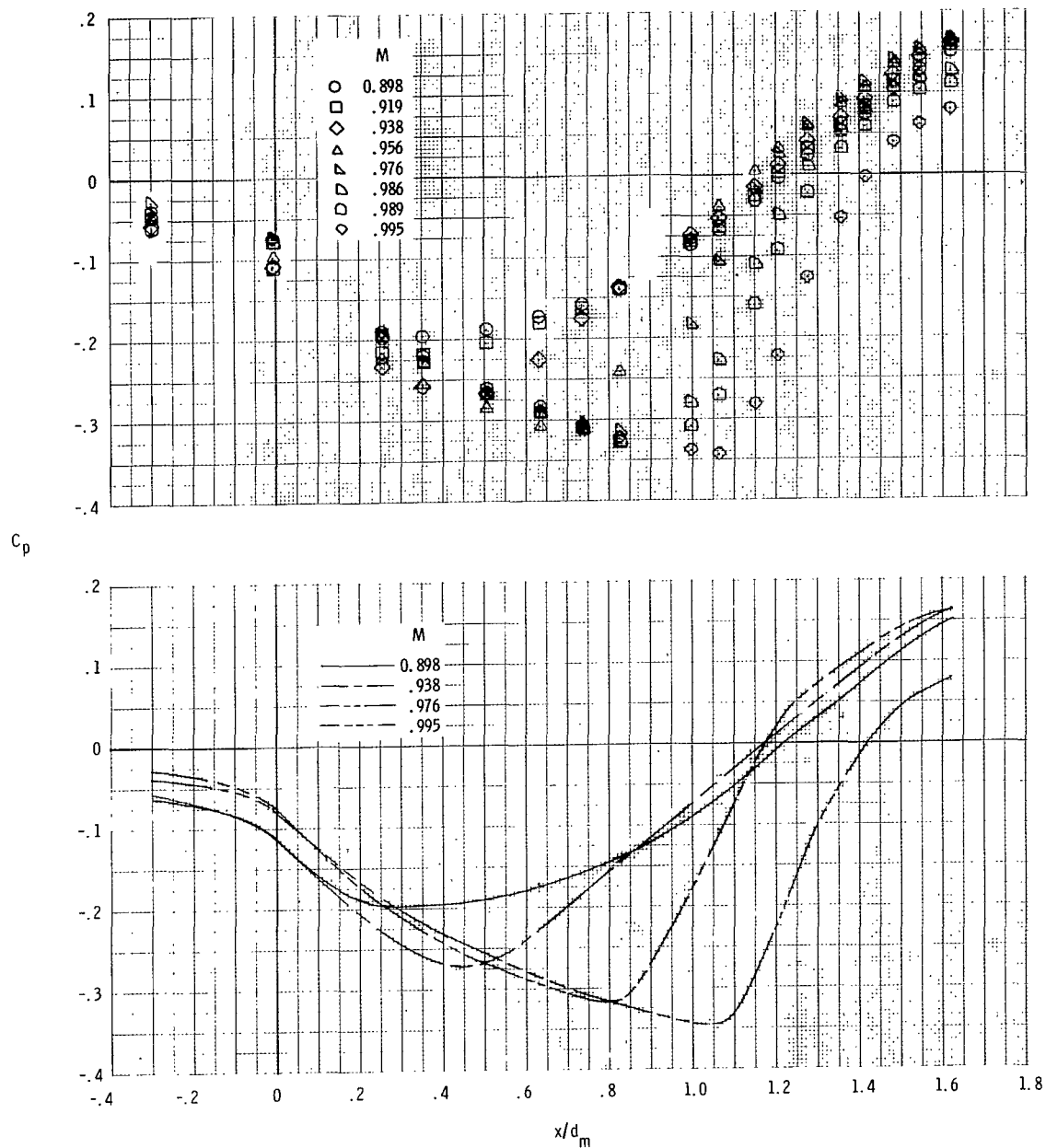
(d)  $M = 1.00$  to  $1.02$ .

Figure 10.- Continued.



(e)  $M = 1.03$  and  $1.05$ .

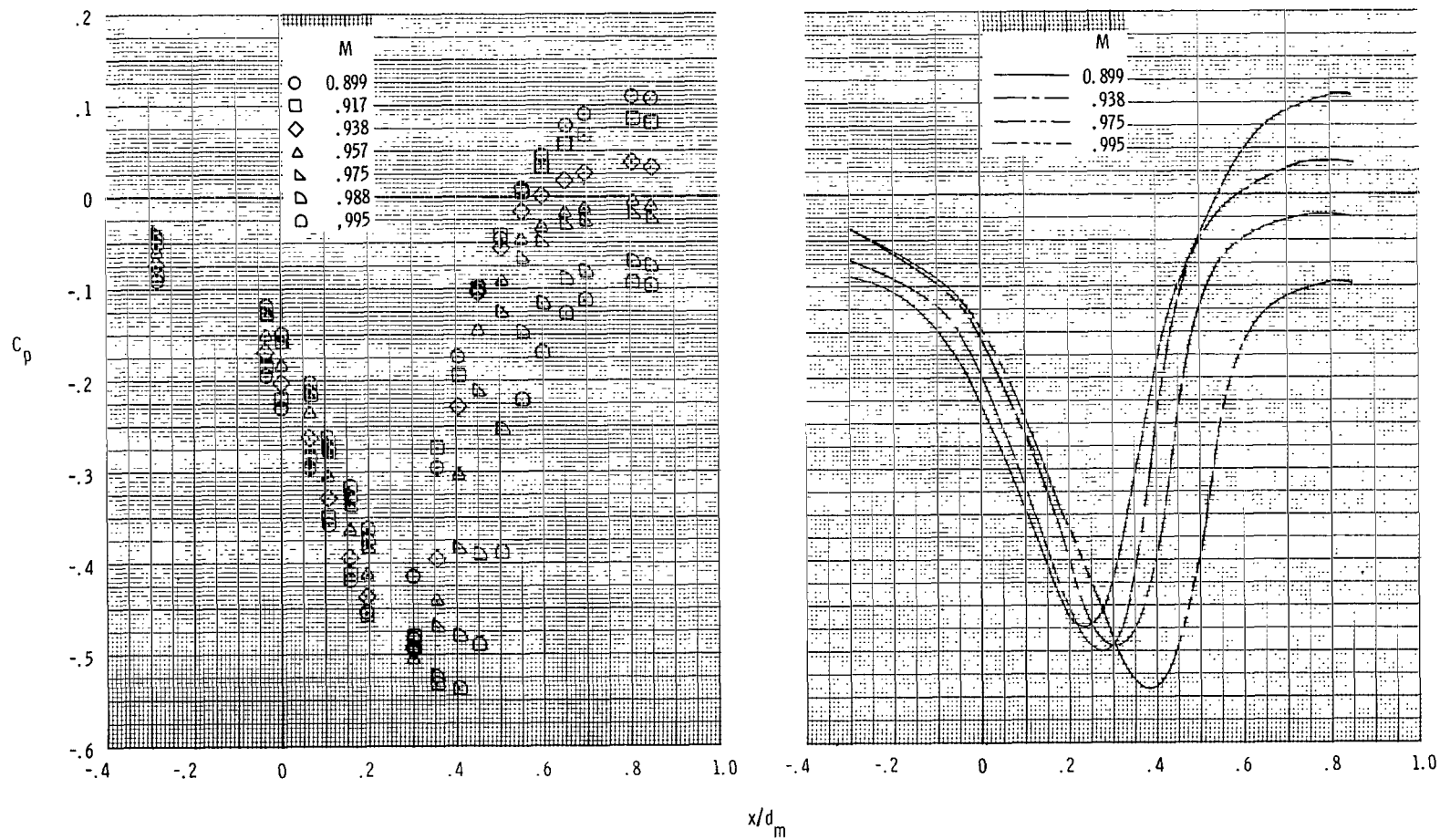
Figure 10.- Concluded.



(a)  $L/d_m = 8.0$ ,  $l/d_m = 1.77$  circular-arc.

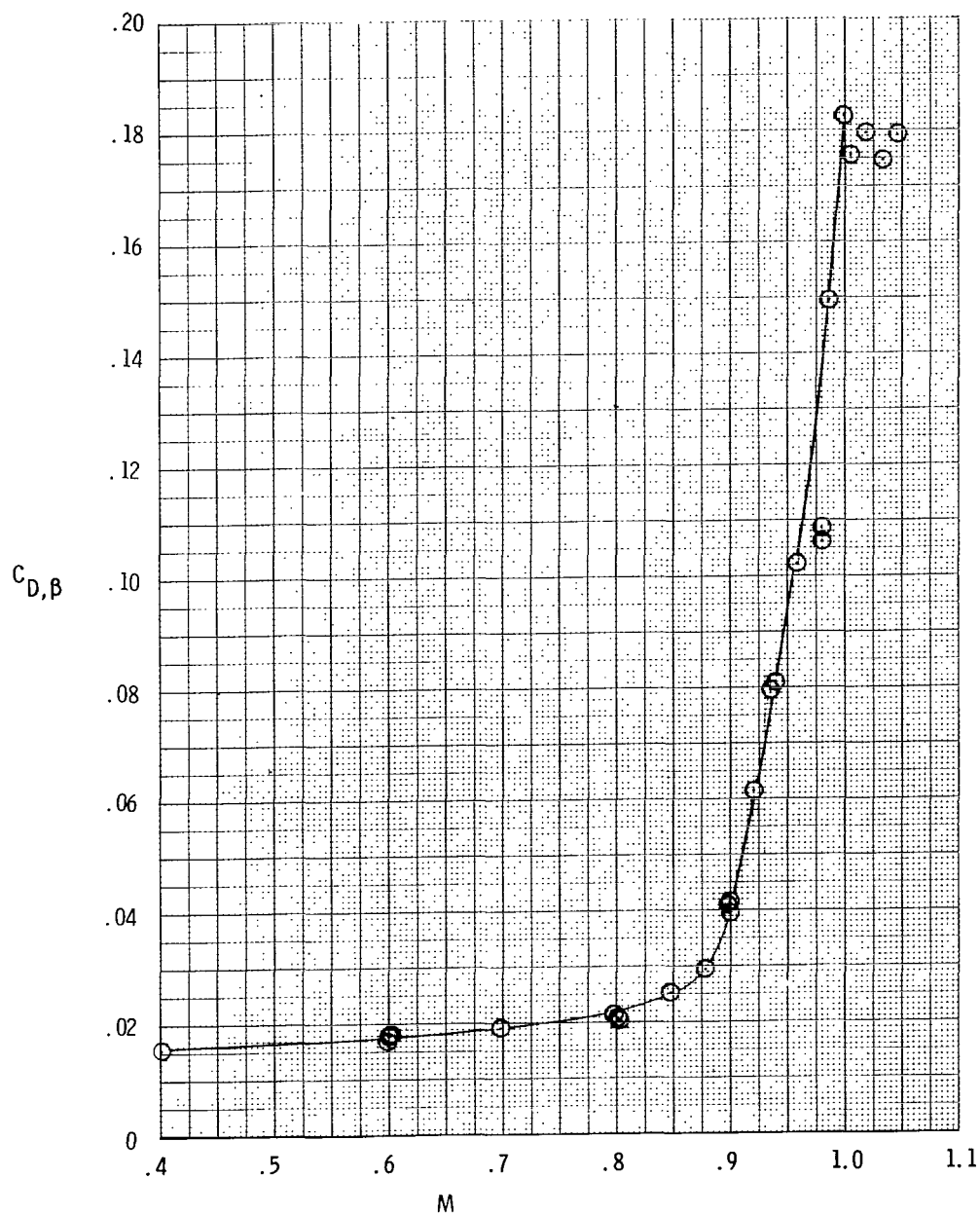
Figure 11.- Typical examples of change in boattail pressure coefficient distributions at Mach numbers between about 0.9 and 1.0.





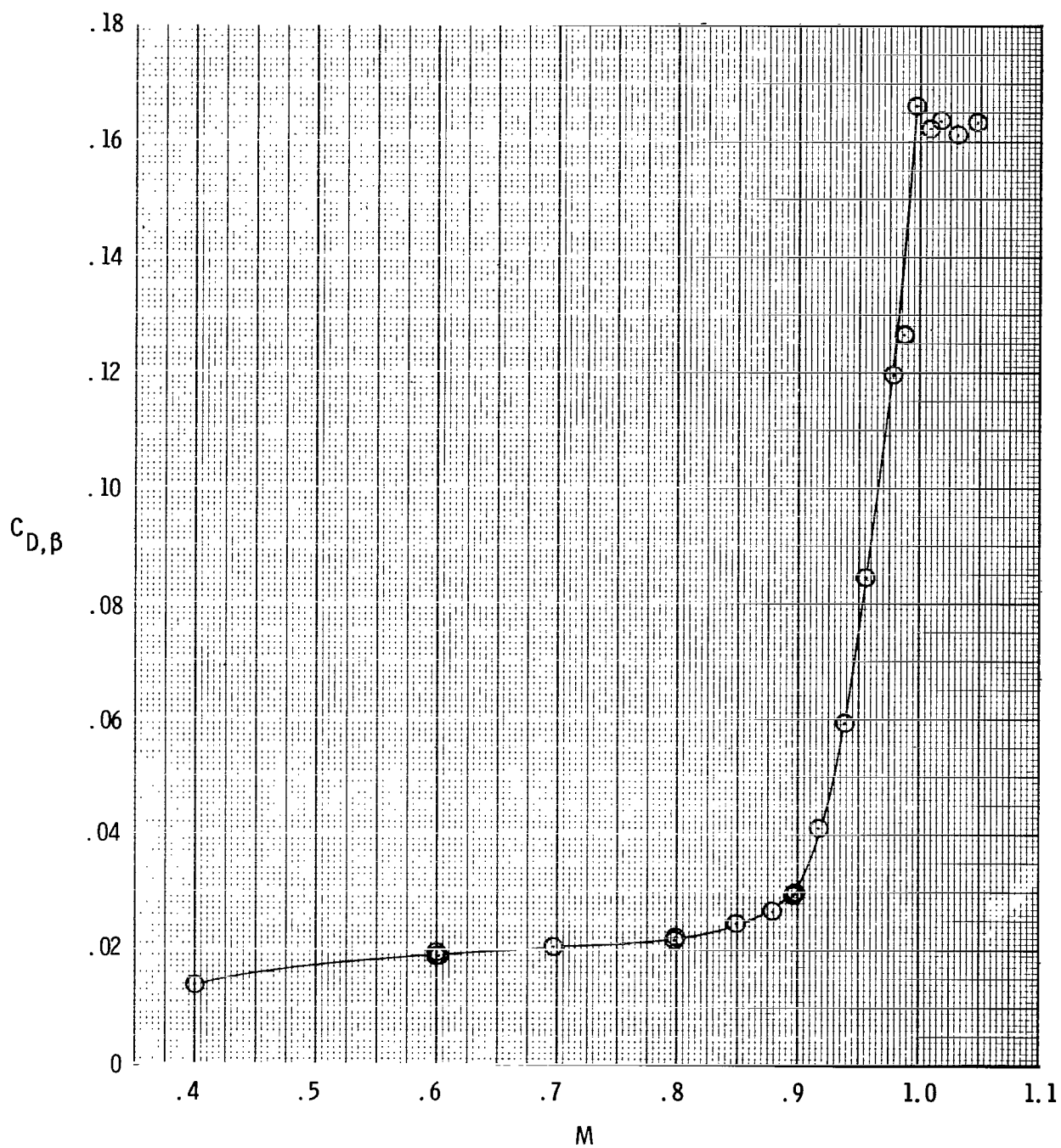
(b)  $L/d_m = 8.0$ ,  $l/d_m = 0.96$  circular-arc-conic.

Figure 11.- Concluded.



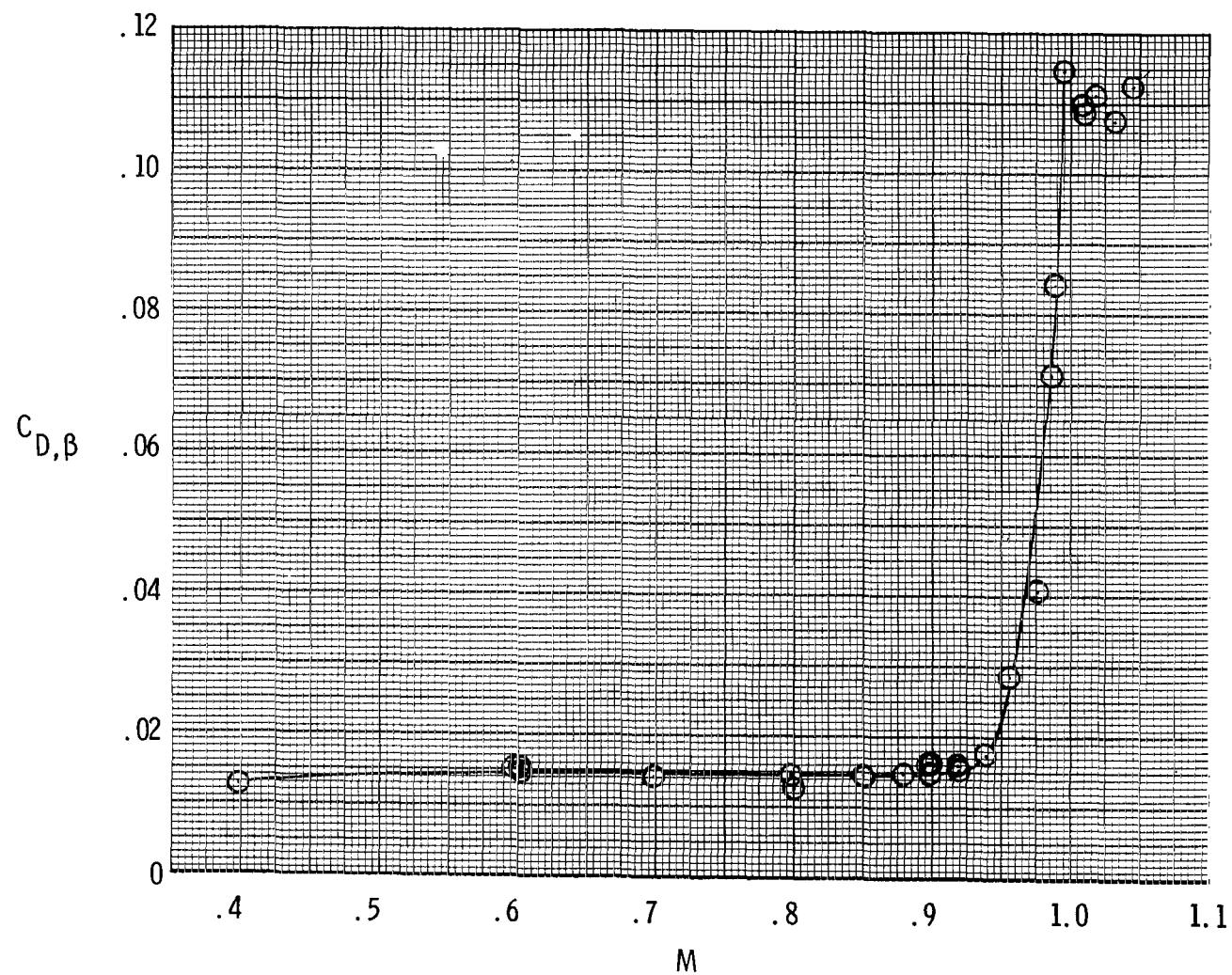
(a)  $L/d_m = 8.0$ ,  $l/d_m = 0.80$  circular-arc.

Figure 12.- Effect of Mach number on boattail pressure drag coefficients for the various configurations.



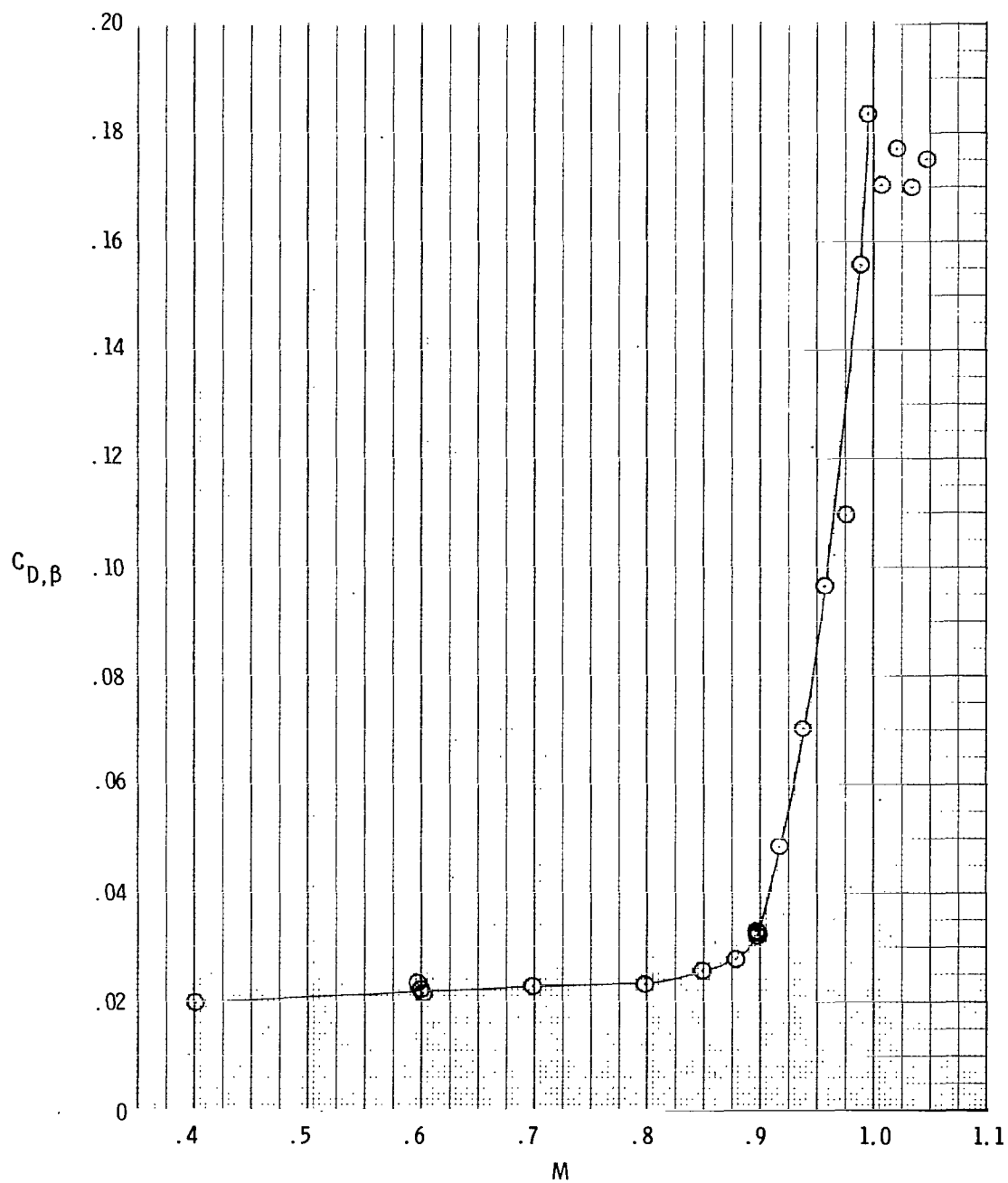
(b)  $L/d_m = 16.0$ ,  $l/d_m = 0.80$  circular-arc.

Figure 12.- Continued.



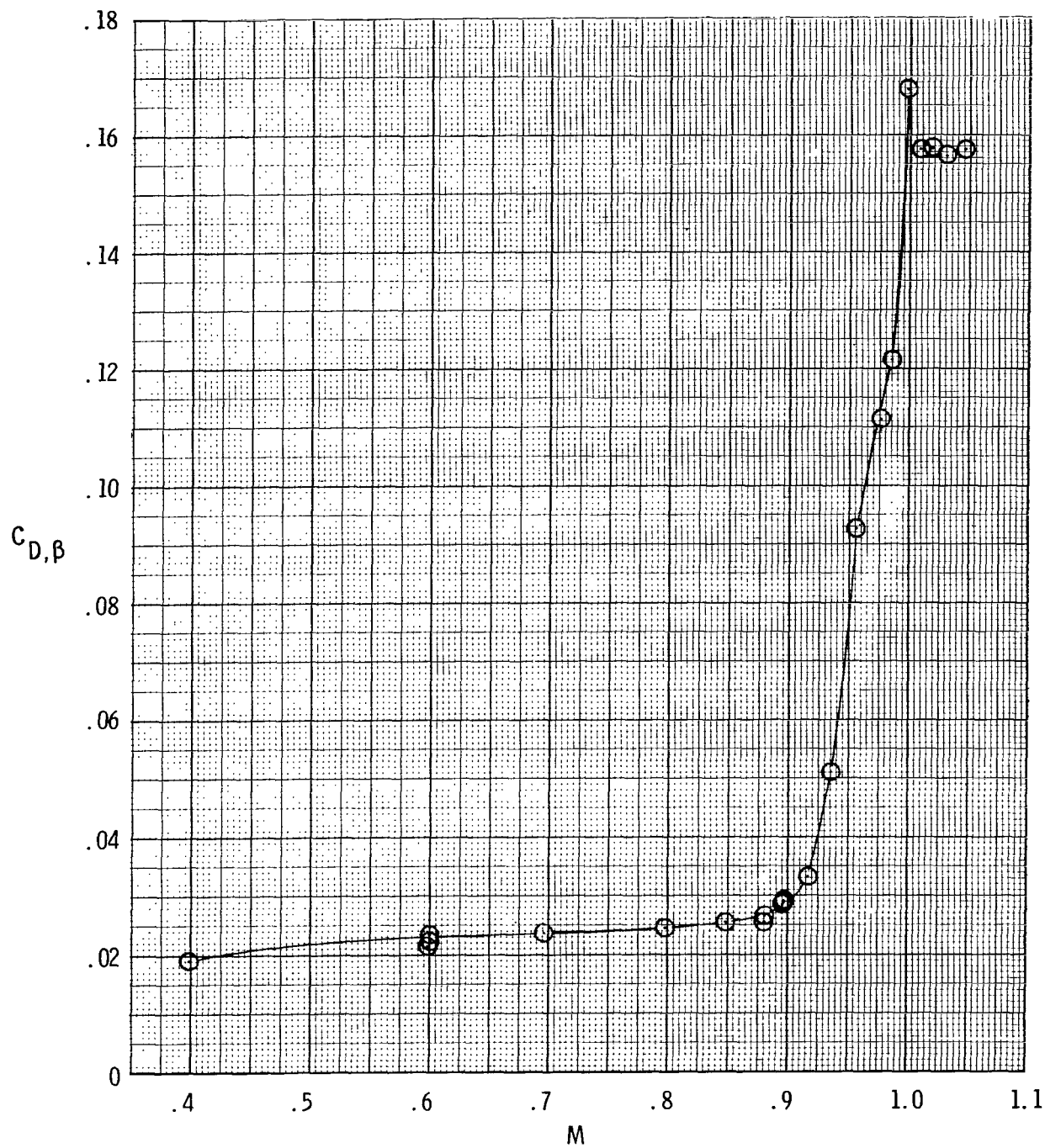
(c)  $L/d_m = 8.0$ ,  $l/d_m = 1.77$  circular-arc.

Figure 12.- Continued.



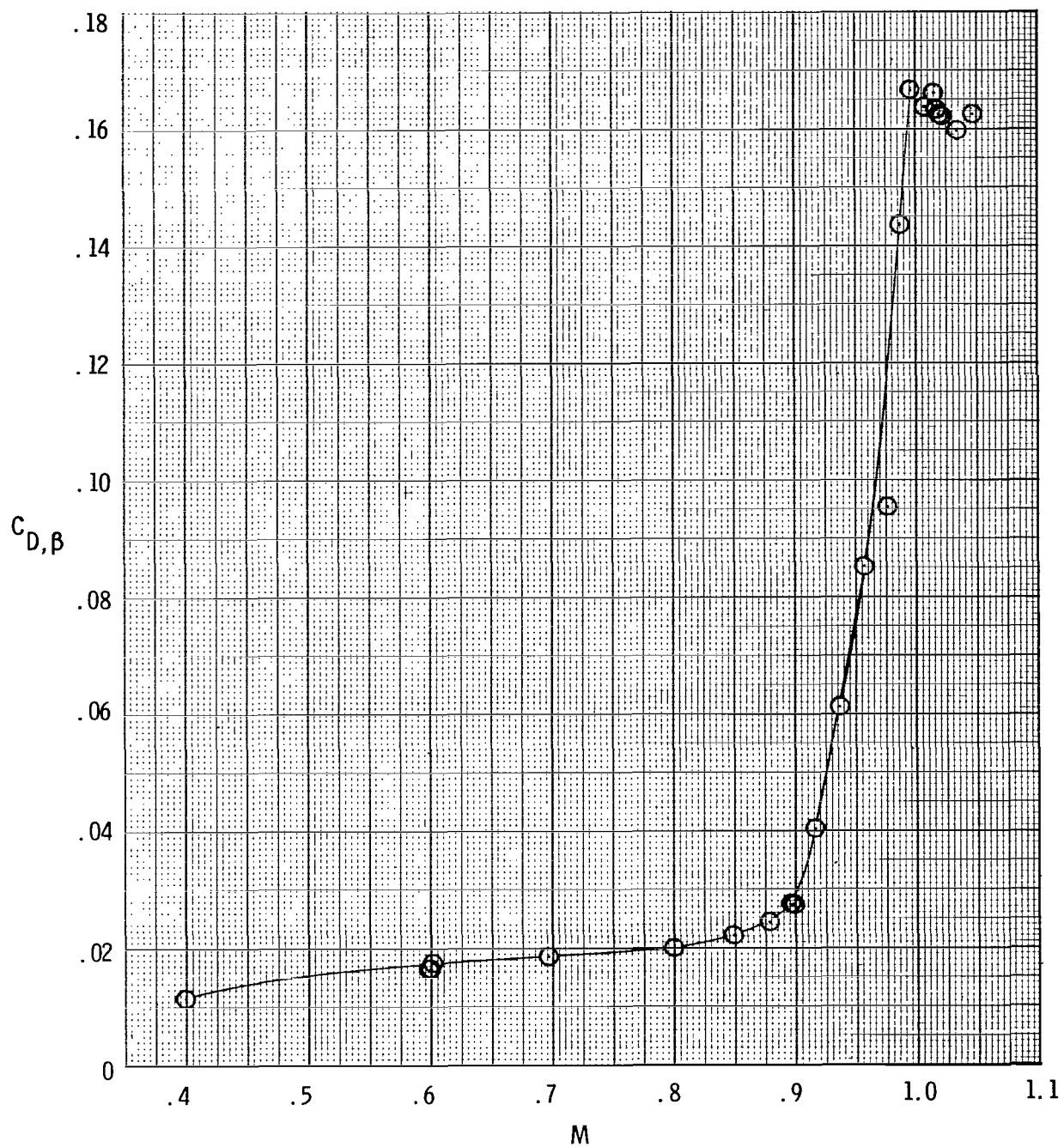
(d)  $L/d_m = 8.0$ ,  $l/d_m = 0.96$  circular-arc-conic.

Figure 12.- Continued.



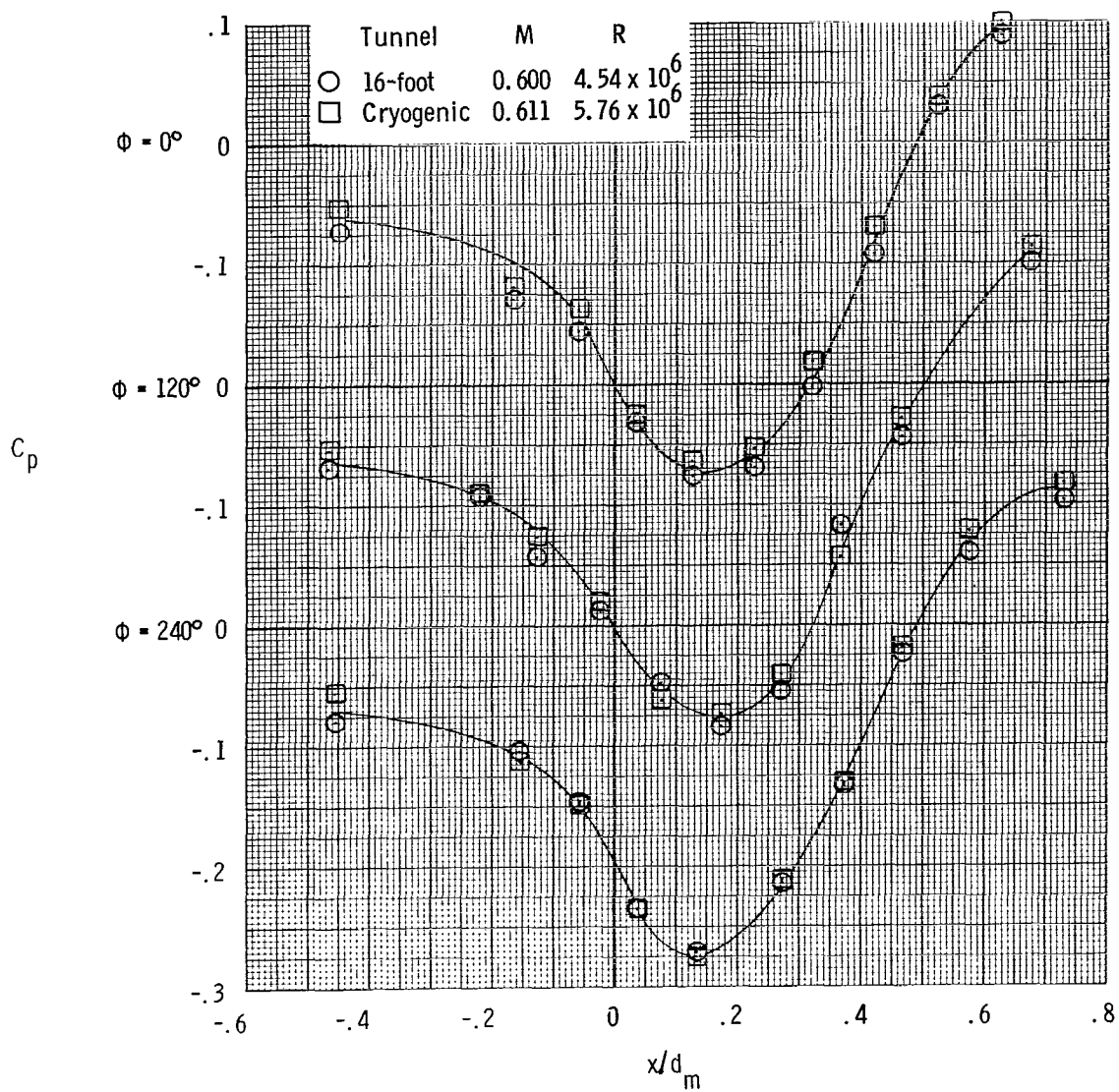
(e)  $L/d_m = 16.0$ ,  $l/d_m = 0.96$  circular-arc-conic.

Figure 12.- Continued.



(f)  $L/d_m = 8.0$ ,  $l/d_m = 0.95$  contoured.

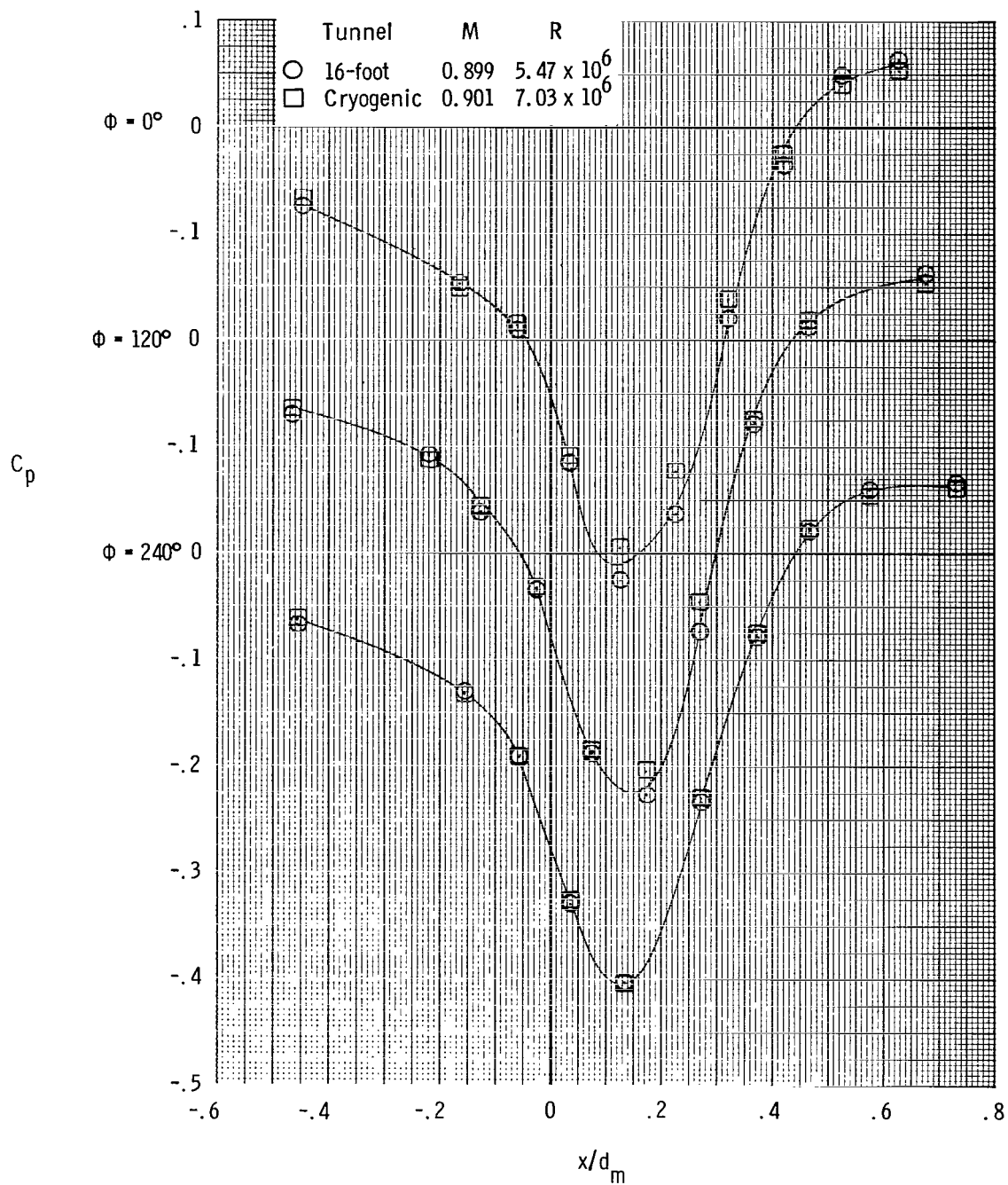
Figure 12.- Concluded.



(a)  $M = 0.6$ .

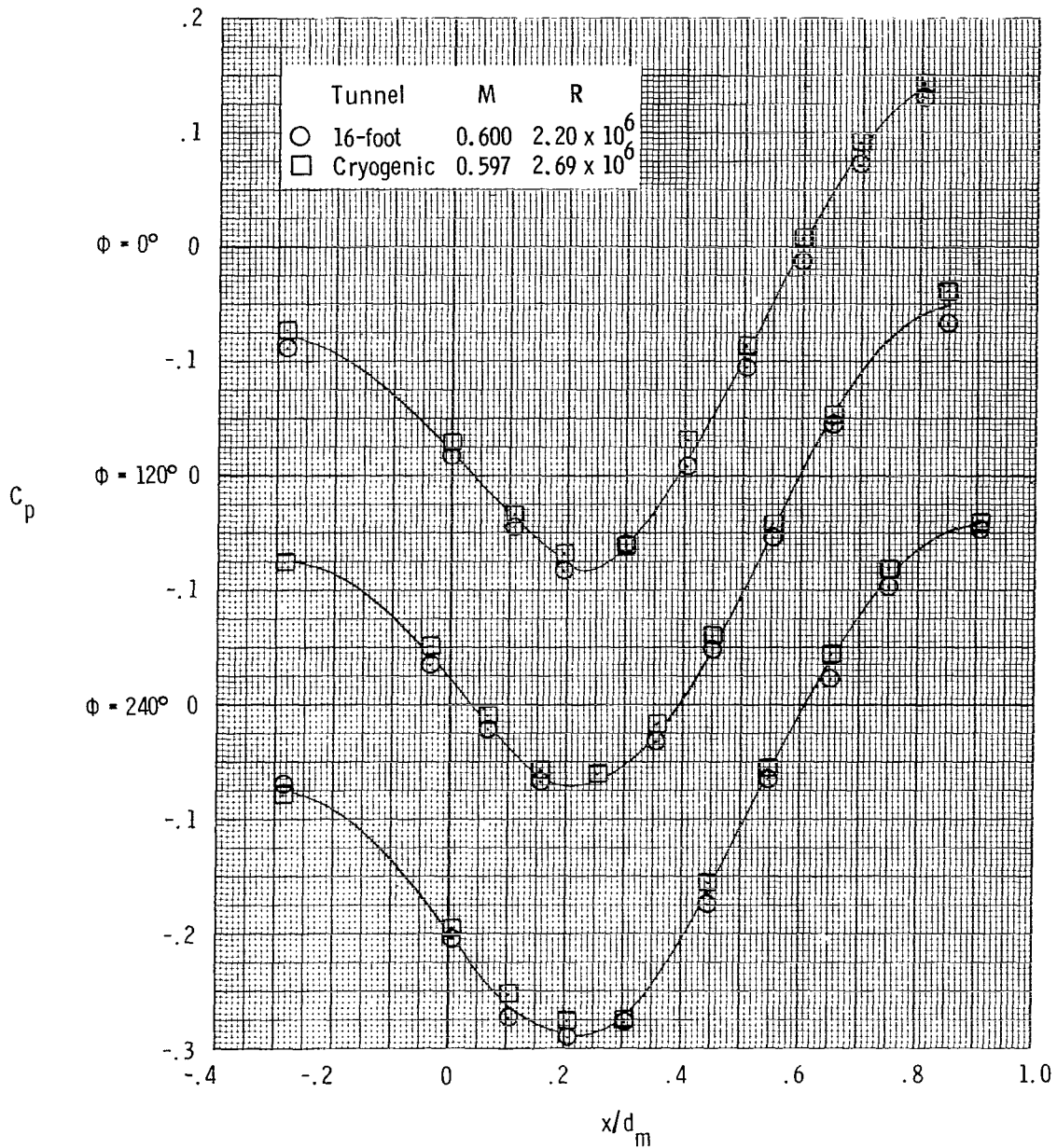
Figure 13.- Boattail pressure coefficient distributions for the  $L/d_m = 16.0$ ,  $l/d_m = 0.80$  circular-arc boattail obtained in both the Langley 16-foot transonic and 0.3-meter transonic cryogenic tunnels at about the same Reynolds number.





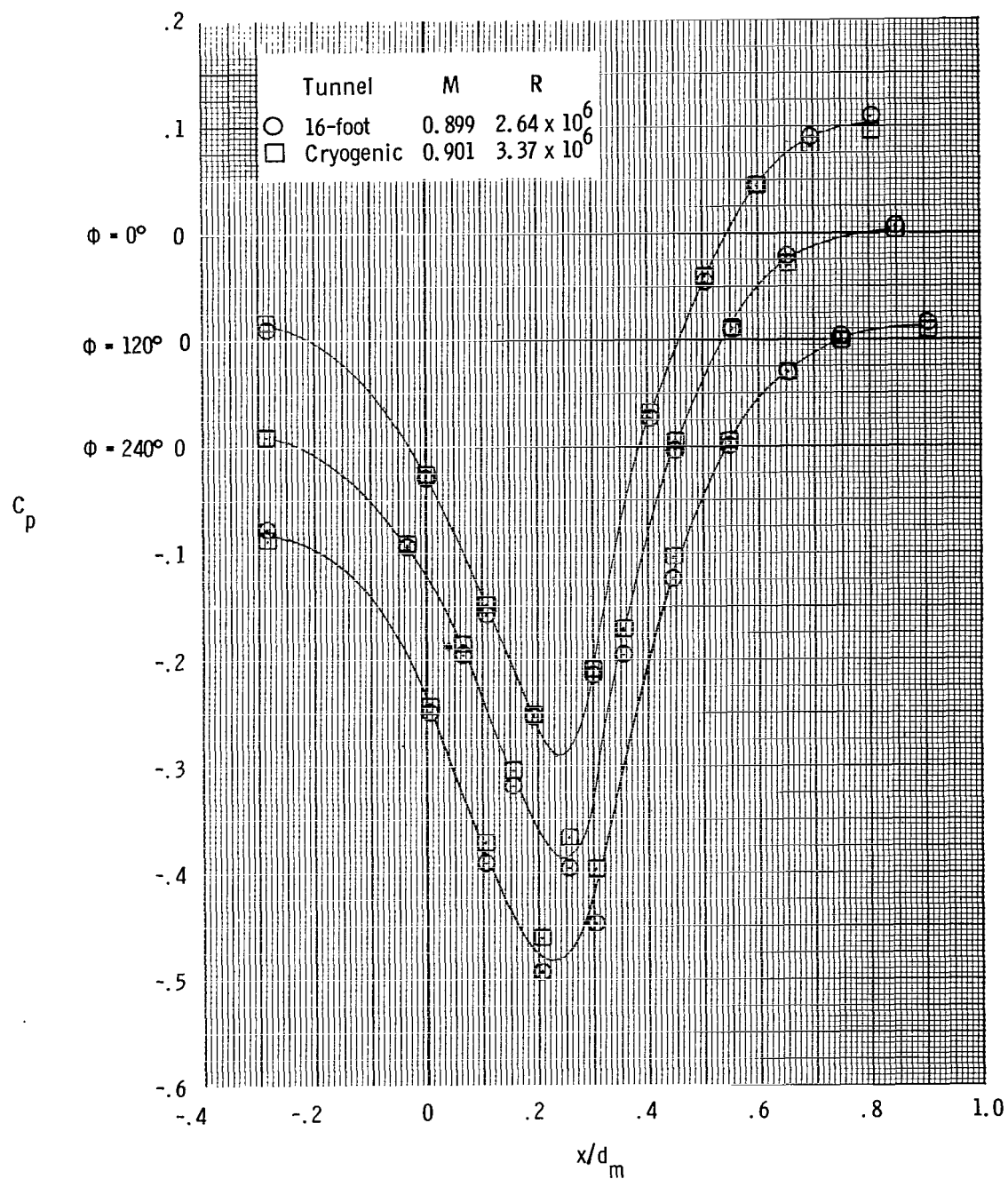
(b)  $M = 0.9$ .

Figure 13.- Concluded.



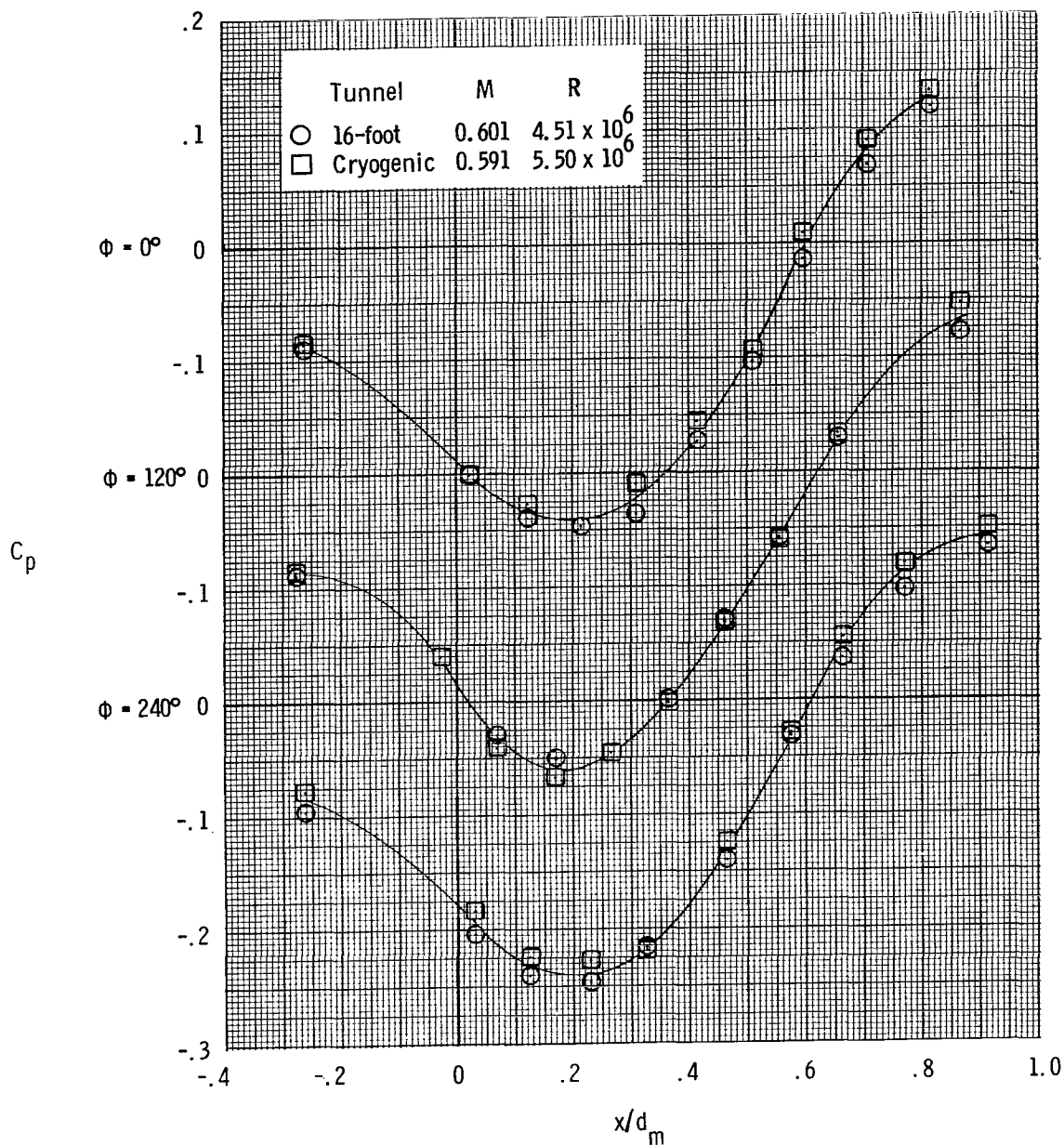
(a)  $M = 0.6$ .

Figure 14.- Boattail pressure coefficient distributions for the  $L/d_m = 8.0$ ,  $l/d_m = 0.96$  circular-arc-conic boattail obtained in both the Langley 16-foot transonic and 0.3-meter transonic cryogenic tunnels at about the same Reynolds number.



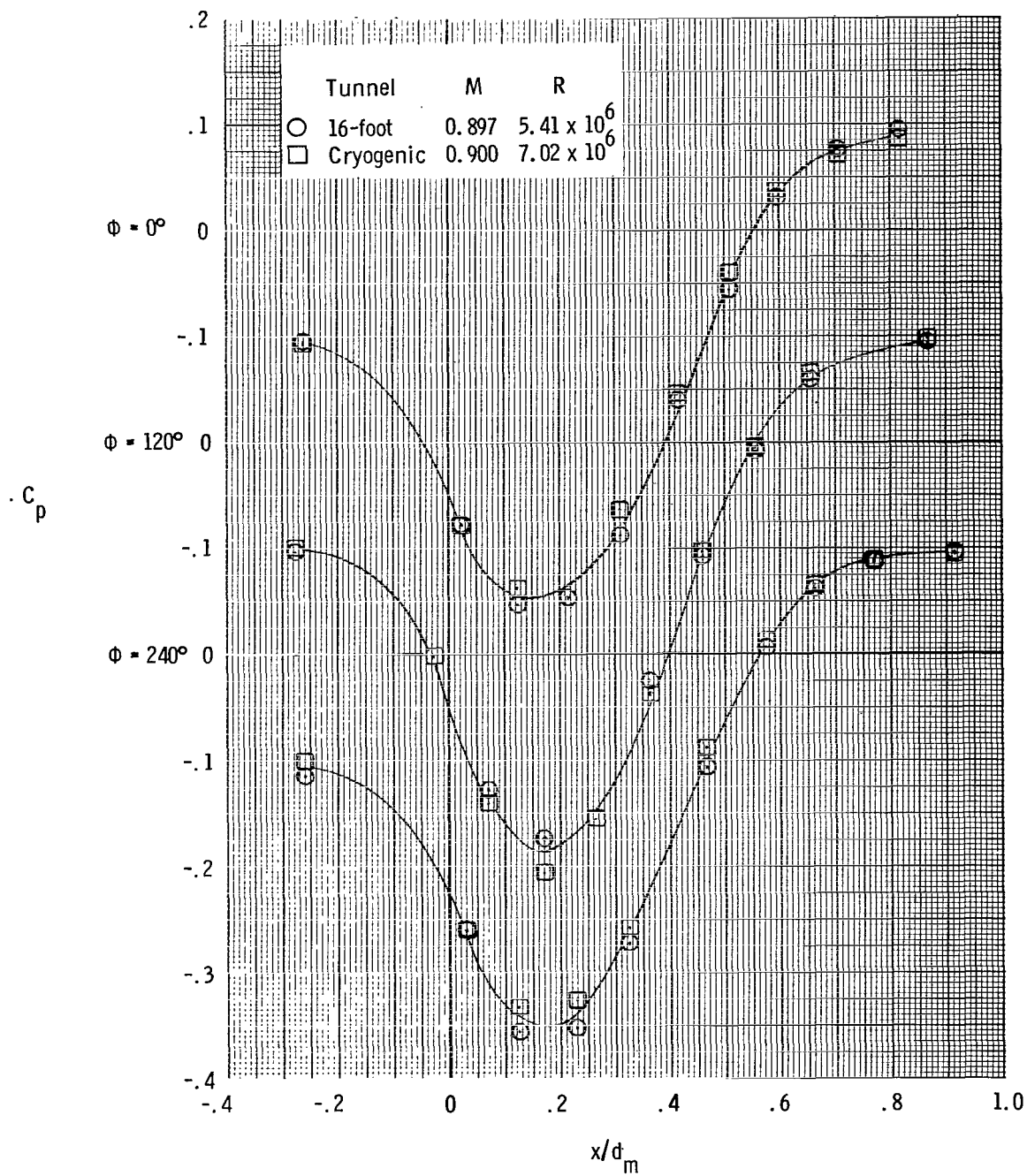
(b)  $M = 0.9$ .

Figure 14.- Concluded.



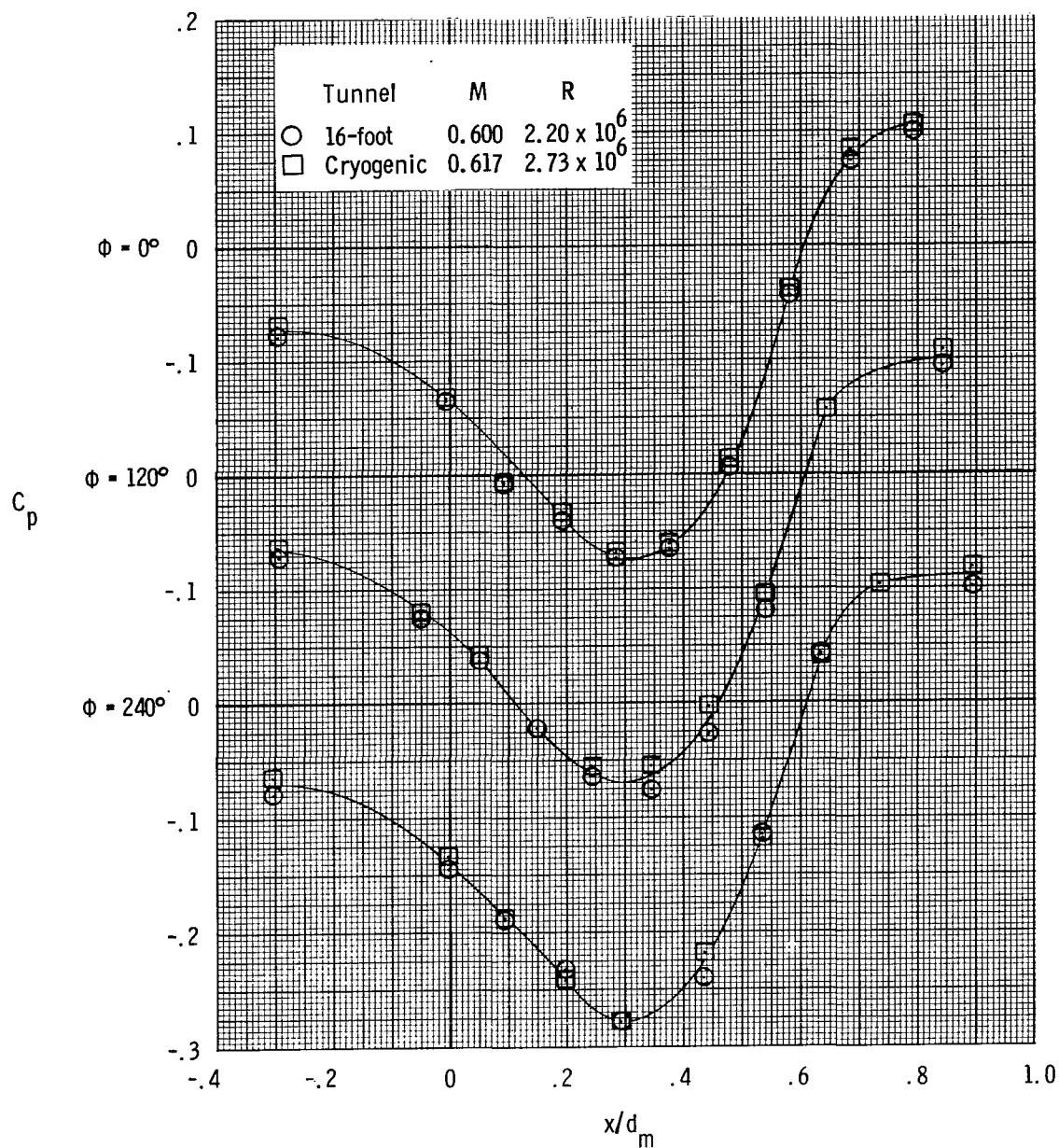
(a)  $M = 0.6$ .

Figure 15.- Boattail pressure coefficient distributions for the  $L/d_m = 16.0$ ,  $l/d_m = 0.96$  circular-arc-conic boattail obtained in both the Langley 16-foot transonic and 0.3-meter transonic cryogenic tunnels at about the same Reynolds number.



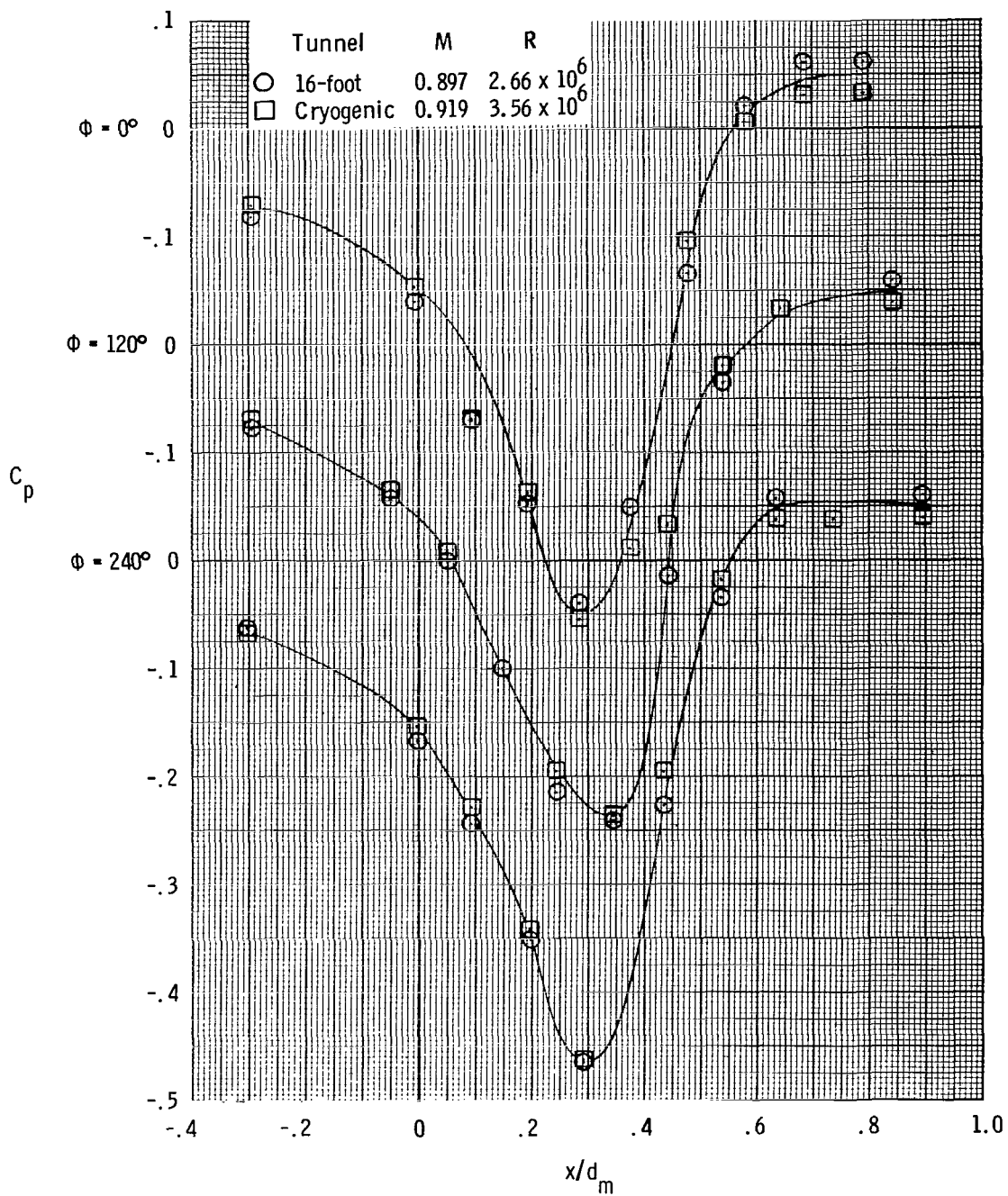
(b)  $M = 0.9$ .

Figure 15.- Concluded.



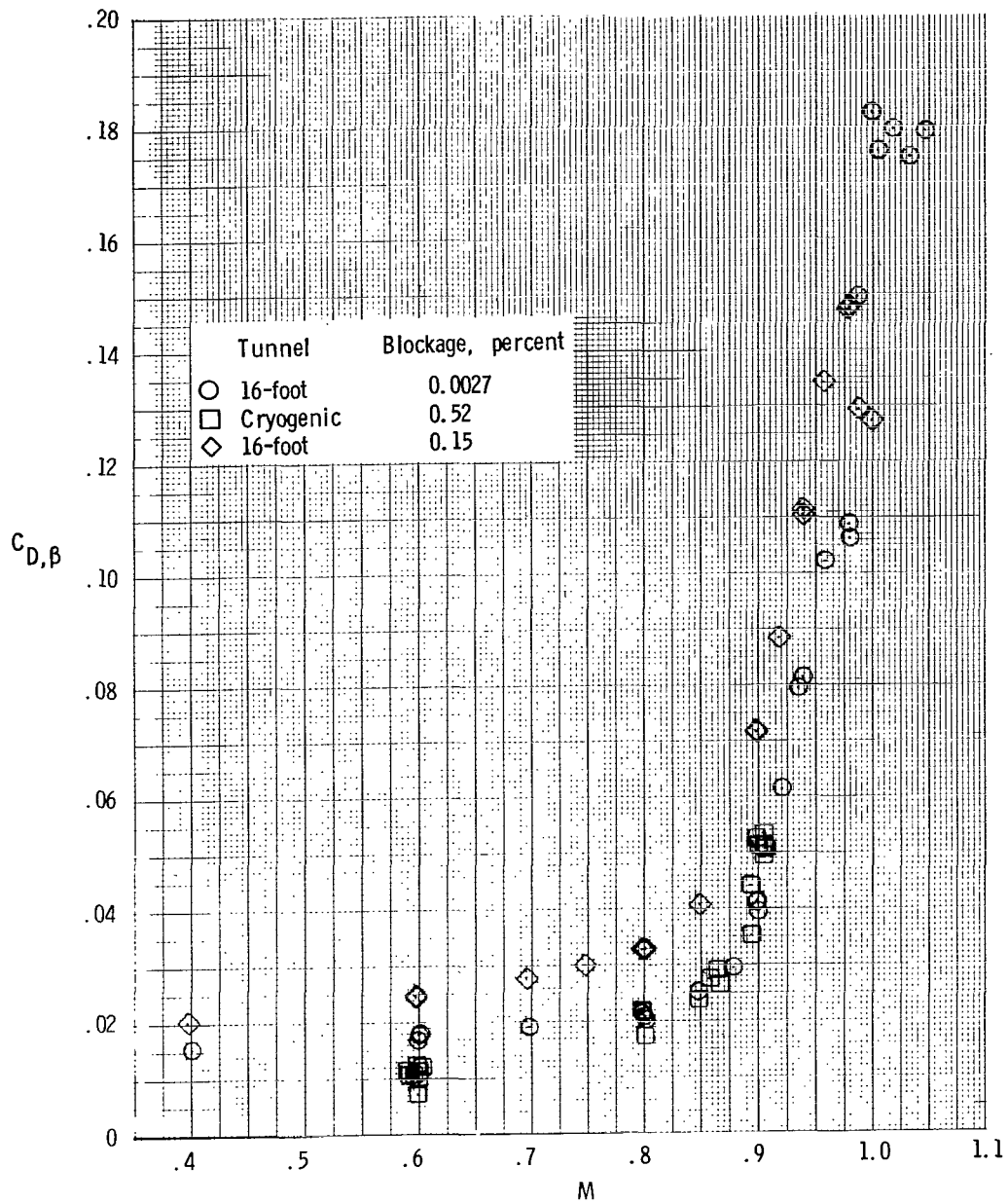
(a)  $M = 0.6$ .

Figure 16.- Boattail pressure coefficient distributions for the  $L/d_m = 8.0$ ,  $l/d_m = 0.95$  contoured boattail obtained in both the Langley 16-foot transonic and 0.3-meter transonic cryogenic tunnels at about the same Reynolds number.



(b)  $M = 0.9$ .

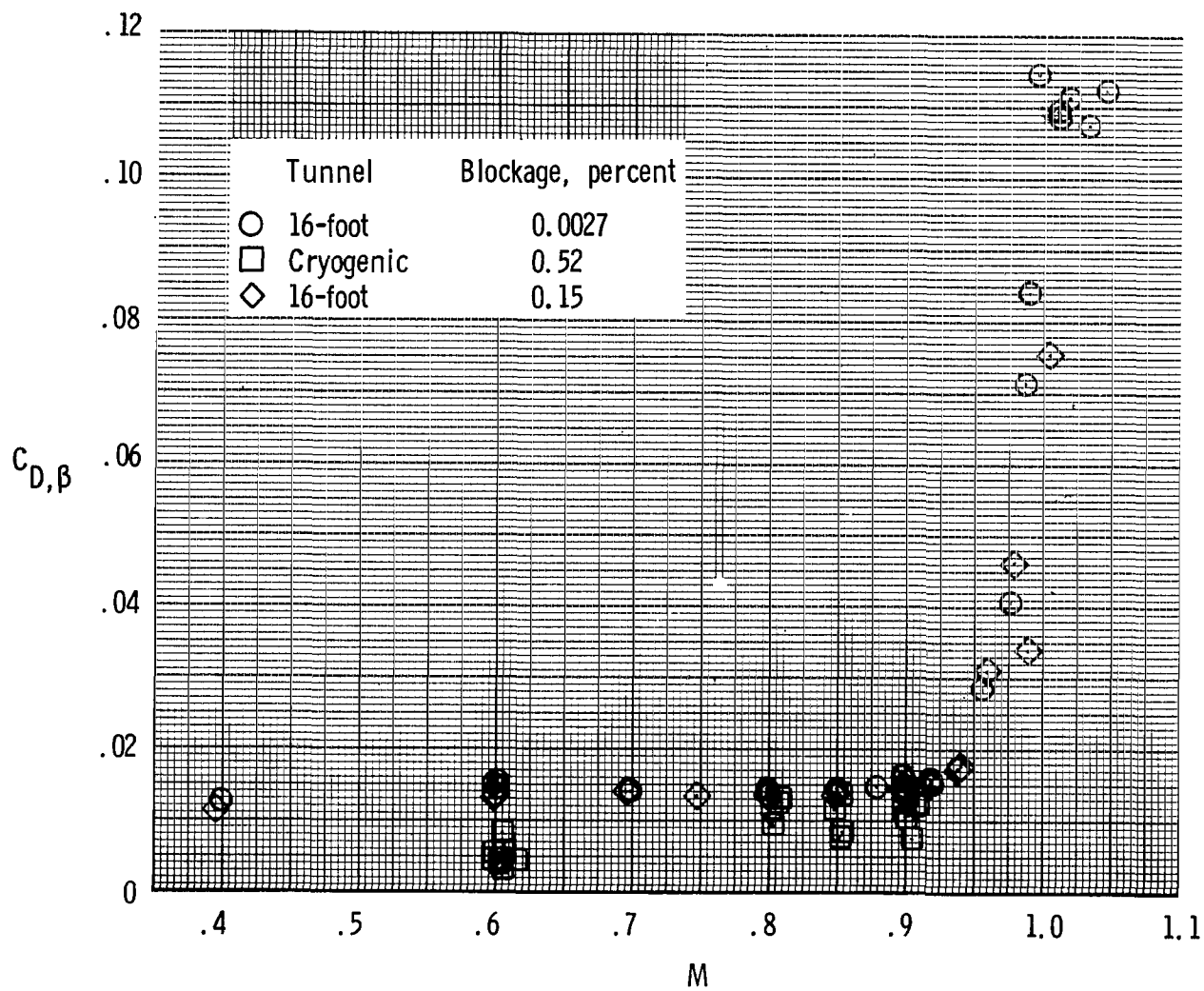
Figure 16.- Concluded.



(a)  $L/d_m = 8.0$ ,  $l/d_m = 0.80$  circular-arc.

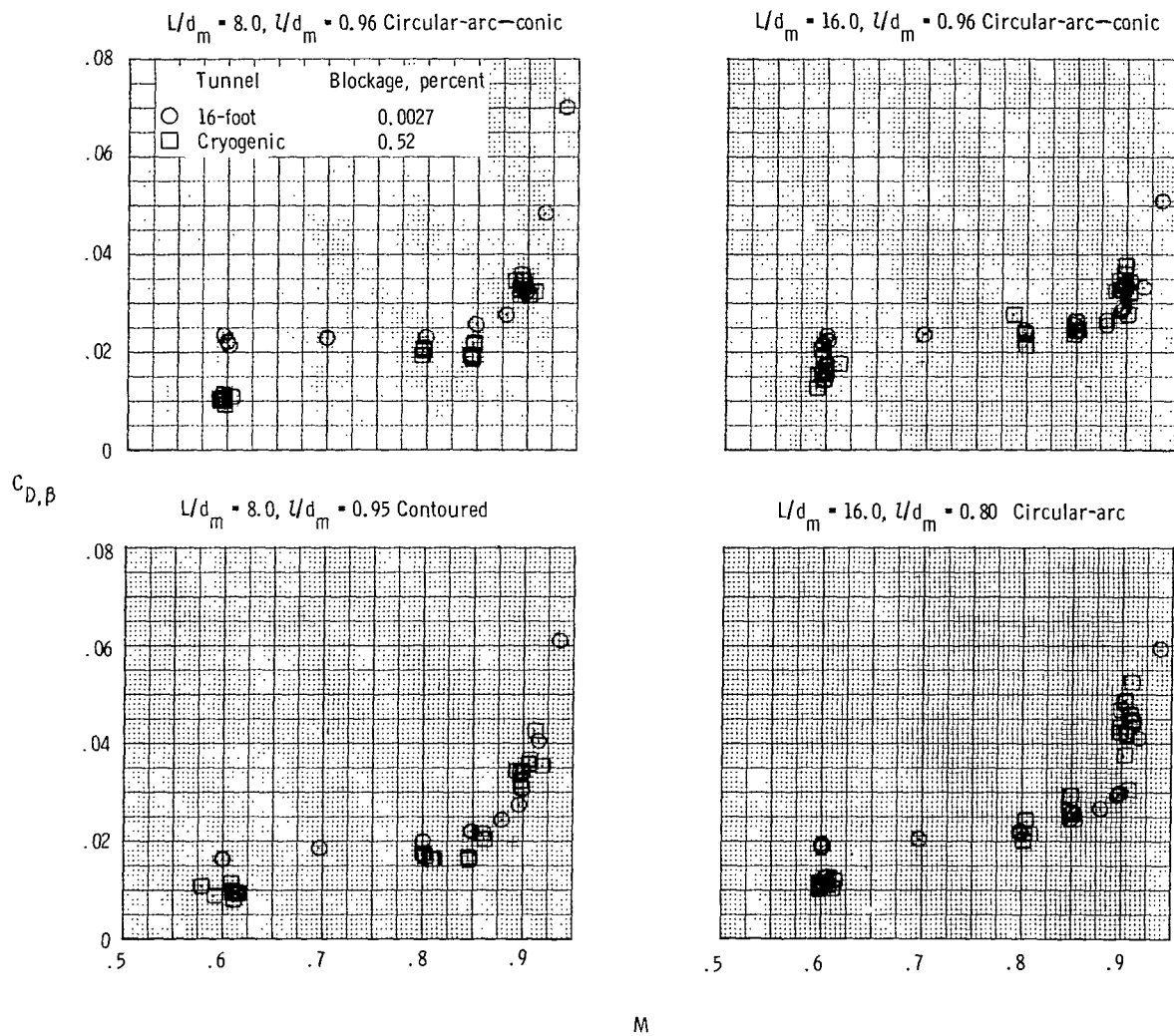
Figure 17.- Boattail pressure drag coefficients for the various boattails obtained in both the Langley 16-foot transonic and 0.3-meter transonic cryogenic tunnels.





(b)  $L/d_m = 8.0$ ,  $l/d_m = 1.77$  circular-arc.

Figure 17.- Continued.



(c) Other configurations.

Figure 17.- Concluded.



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